

Breaking the diffraction limit: Analysis of diode lasers by nearfield scanning optical microscopy (NSOM)

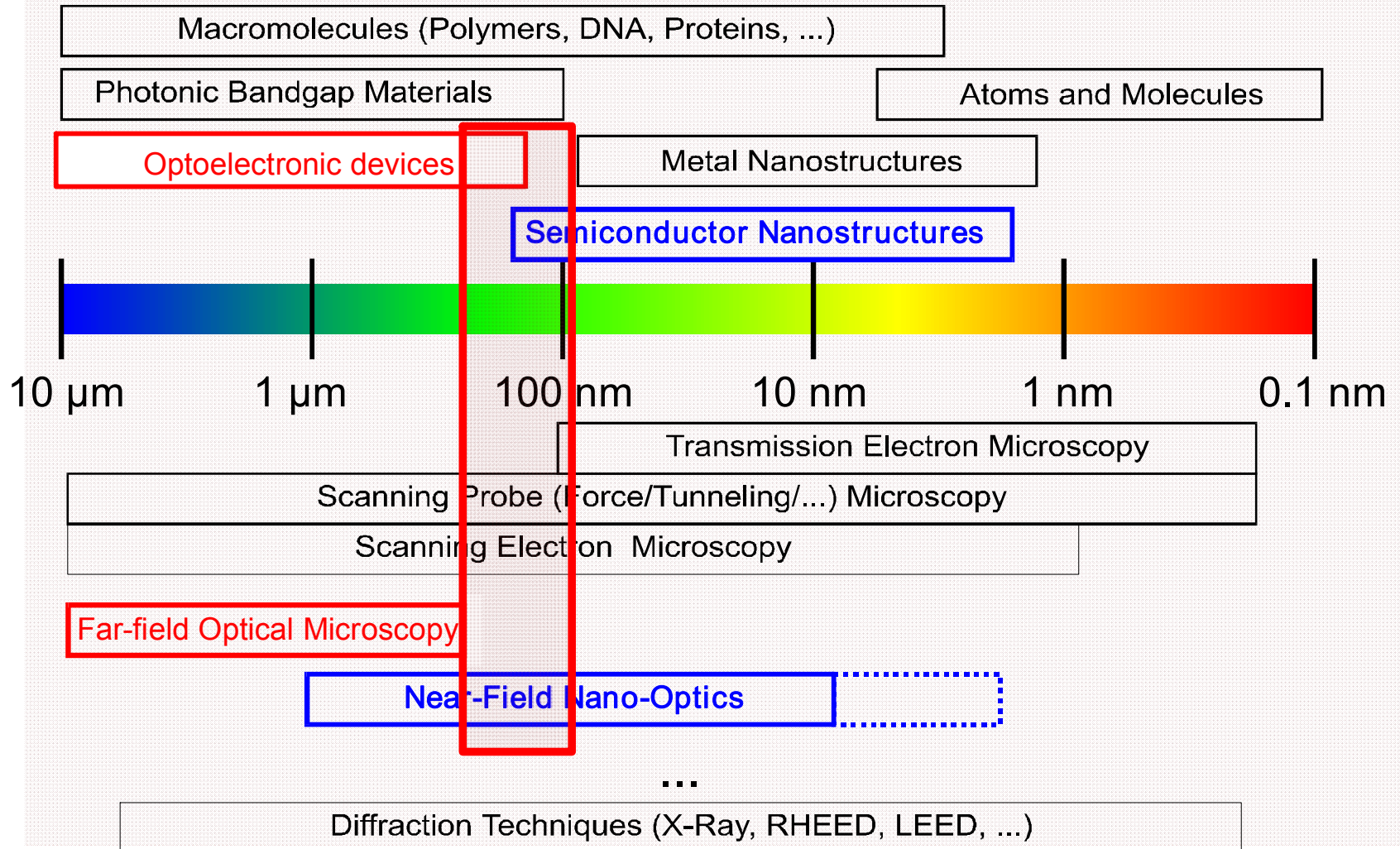
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Ammerländer Heerstraße 114-118, D-26129 Oldenburg, Germany
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1. Introduction
 - NSOM = Nearfield Scanning Optical Microscopy
 - Principles and opportunities
 - Spatial resolution
2. Methodology
 - Laser emission, spontaneous emission and PL
3. NOBIC = Nearfield Optical Beam Induced Current
4. Experimental setups and equipment
5. Analysis of waveguides and determination of mode profiles
6. VCSEL
7. Surface recombination velocity at facets
8. Defect creation during device operation
9. Summary
10. Acknowledgement

Nanostructures and Nanoanalytical Tools

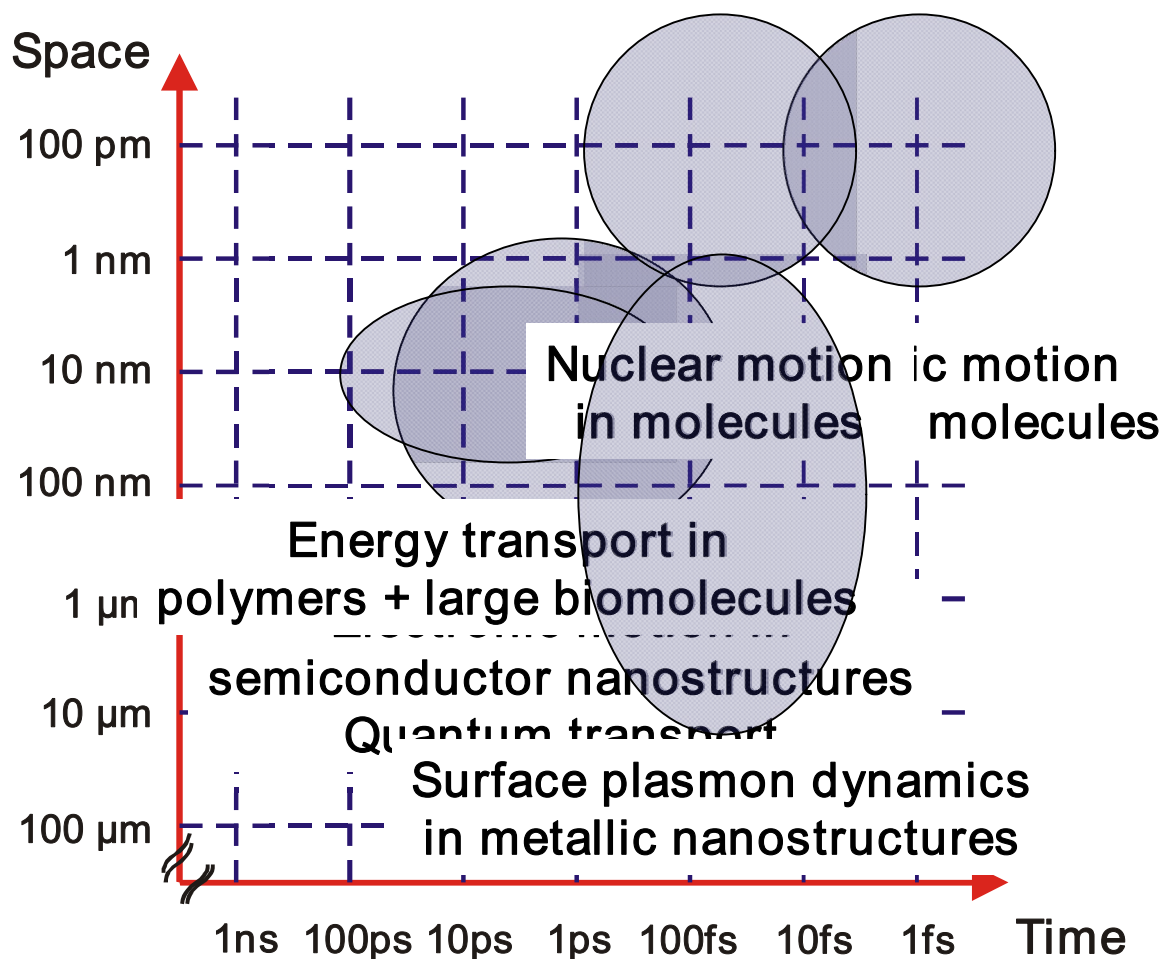


Scanning probe microscopy

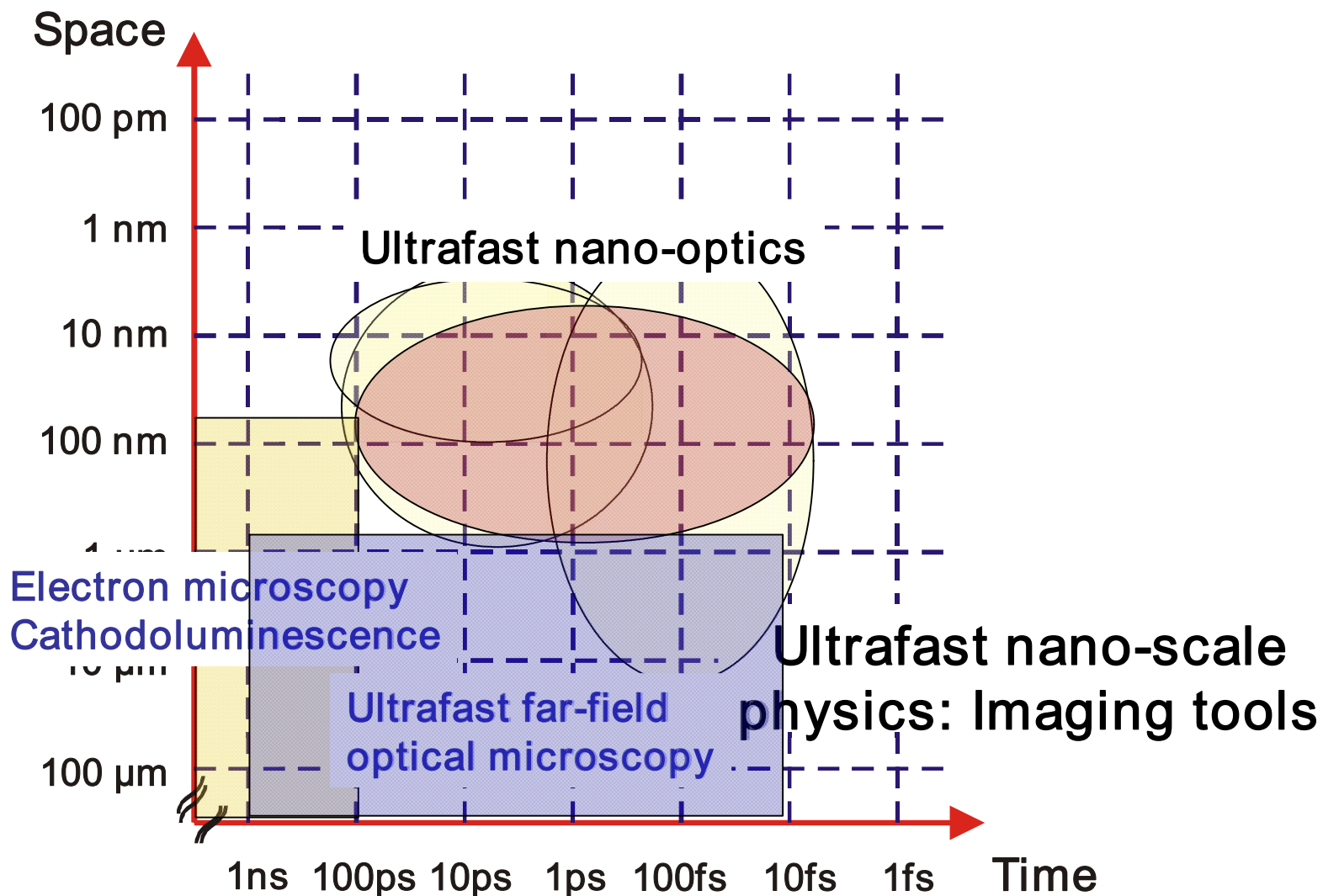


Compare: If 1 atom had the size of an orange,
the cantilever would be 100 km long

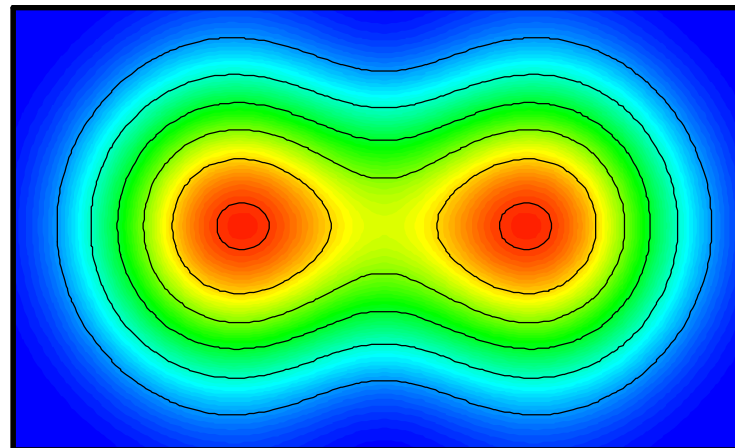
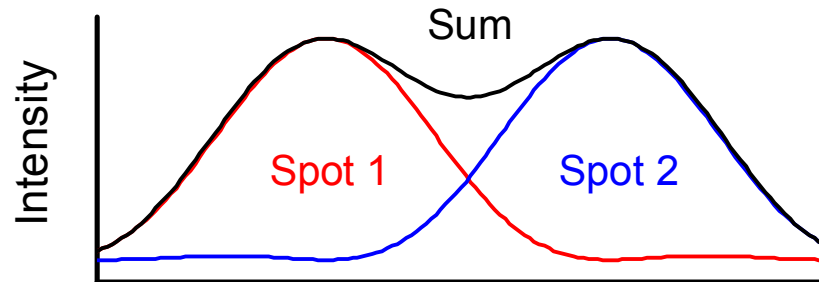
Disc players for atoms



Physical phenomena happen on ultrashort length and time scales



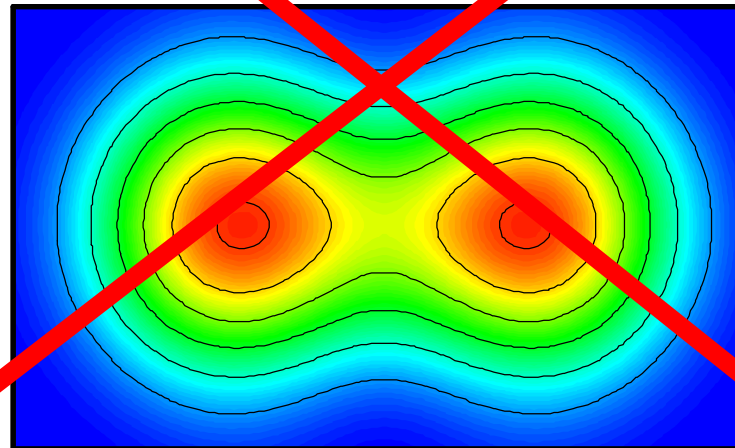
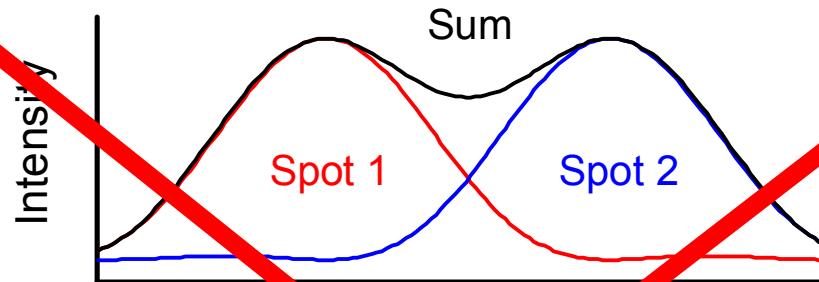
Rayleigh limit: $\Delta x \approx 0.61 \lambda / \text{N.A.}$



The resolution in optical microscopy is limited by the wavelength of light

Rayleigh limit:

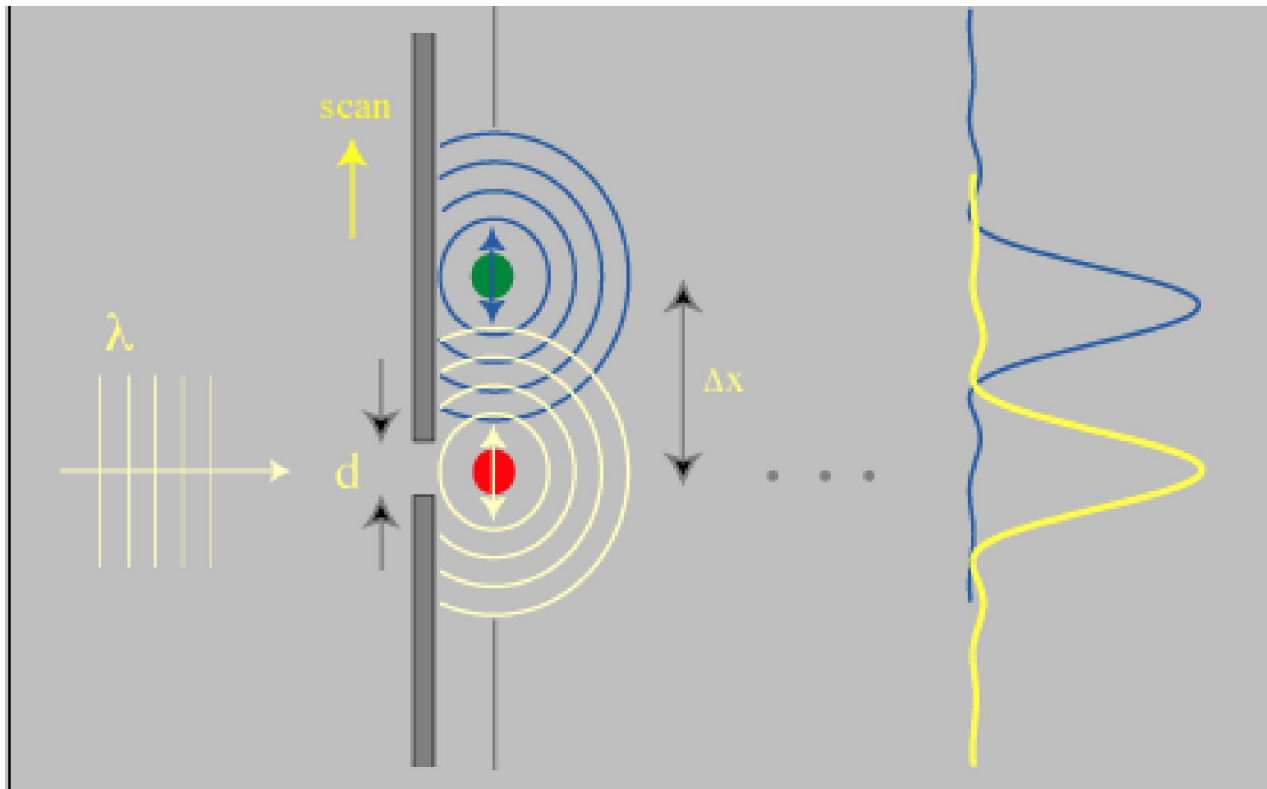
$$\Delta x \approx 0.61 \lambda / \text{N.A.}$$



The resolution in optical microscopy is limited by the wavelength of light

Breaking the resolution limit in microscopy

Near-field scanning optical microscopy



Diffraction-unlimited resolution

Spatial resolution in near-field microscopy

Electromagnetic-field distribution: superposition of monochromatic plane waves:

$$E = E_0 \exp(i\vec{k}_0 \vec{r}) \exp(-i\omega t) = E_0 \exp(i(k_x x + k_y y)) \exp(ik_z z) \exp(-i\omega t)$$

$$k_0 = |\vec{k}_0| = \sqrt{k_x^2 + k_y^2 + k_z^2}$$

for given k_x, k_y : two solutions:

- (a) propagating waves: $k_{lat} = \sqrt{k_x^2 + k_y^2} < k_0 \quad k_z \text{ real.}$
- (b) evanescent waves: $k_{lat} = \sqrt{k_x^2 + k_y^2} > k_0 \quad k_z \text{ imaginary.}$

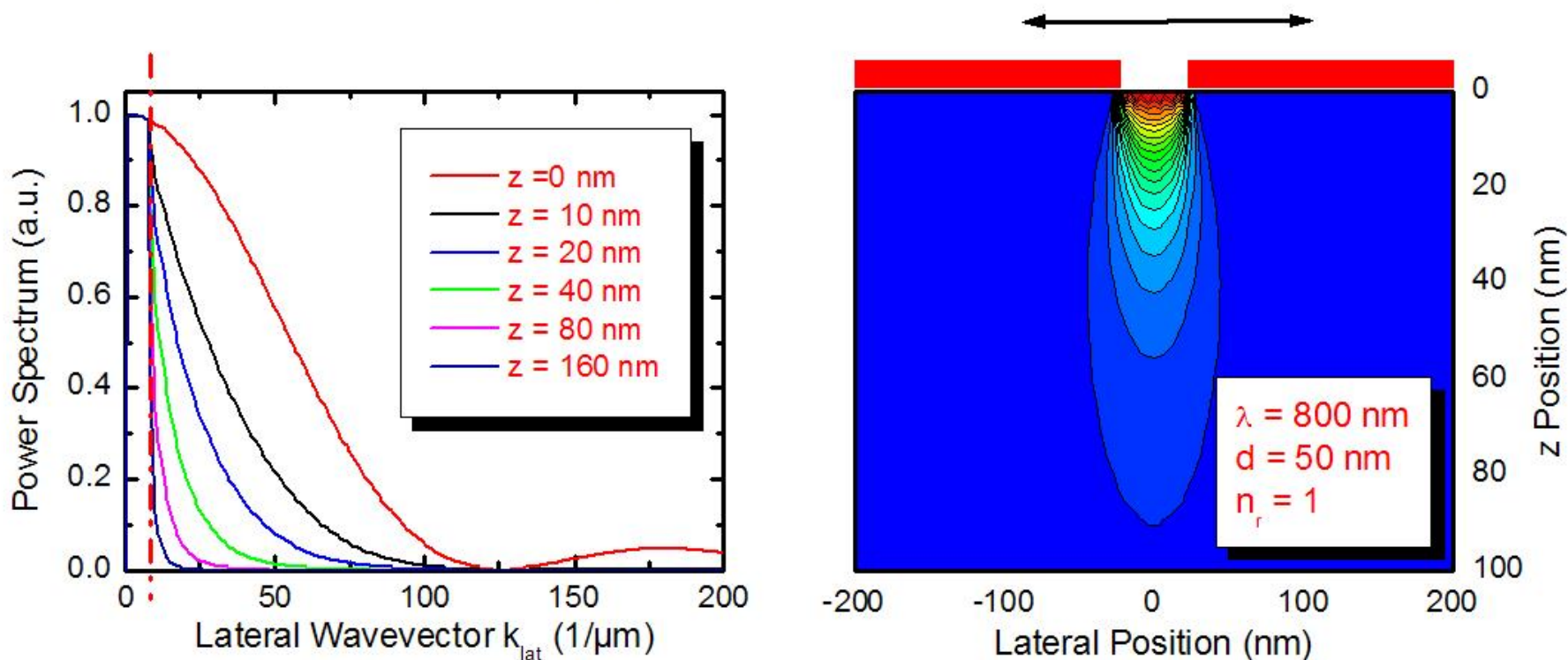
Consequence: Intensity of evanescent waves decreases exponentially with increasing z

How to get optical super-resolution ? : $\Delta k_x \Delta x \geq 1$

Use evanescent modes!

Diffraction by 1-dimensinal slit (Kirchhoff Approximation)

Monochromatic plane wave ($k_0 = 2\pi n_r/\lambda$ - $\lambda = 800$ nm) incident on slit (width $d = 50$ nm)

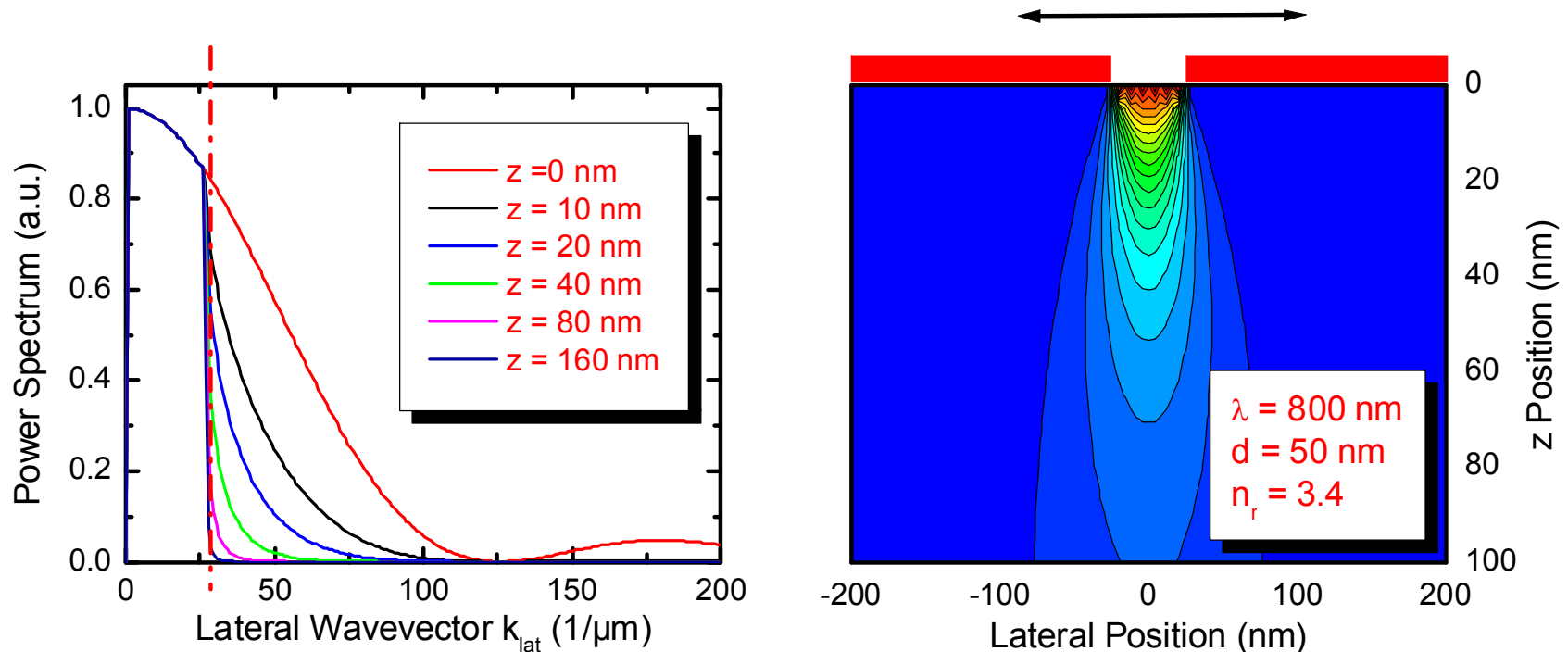


Evanescent waves: $k_{lat} > k_0 = 2\pi n_r/\lambda$ - $k_z^2 = (k_0^2 - k_{lat}^2) < 0$ - k_z imaginary

Propagating waves: $k_{lat} < k_0 = 2\pi n_r/\lambda$ - $k_z^2 = (k_0^2 - k_{lat}^2) > 0$ - k_z real

Diffraction by 1-dimensional slit (Kirchhoff Approximation)

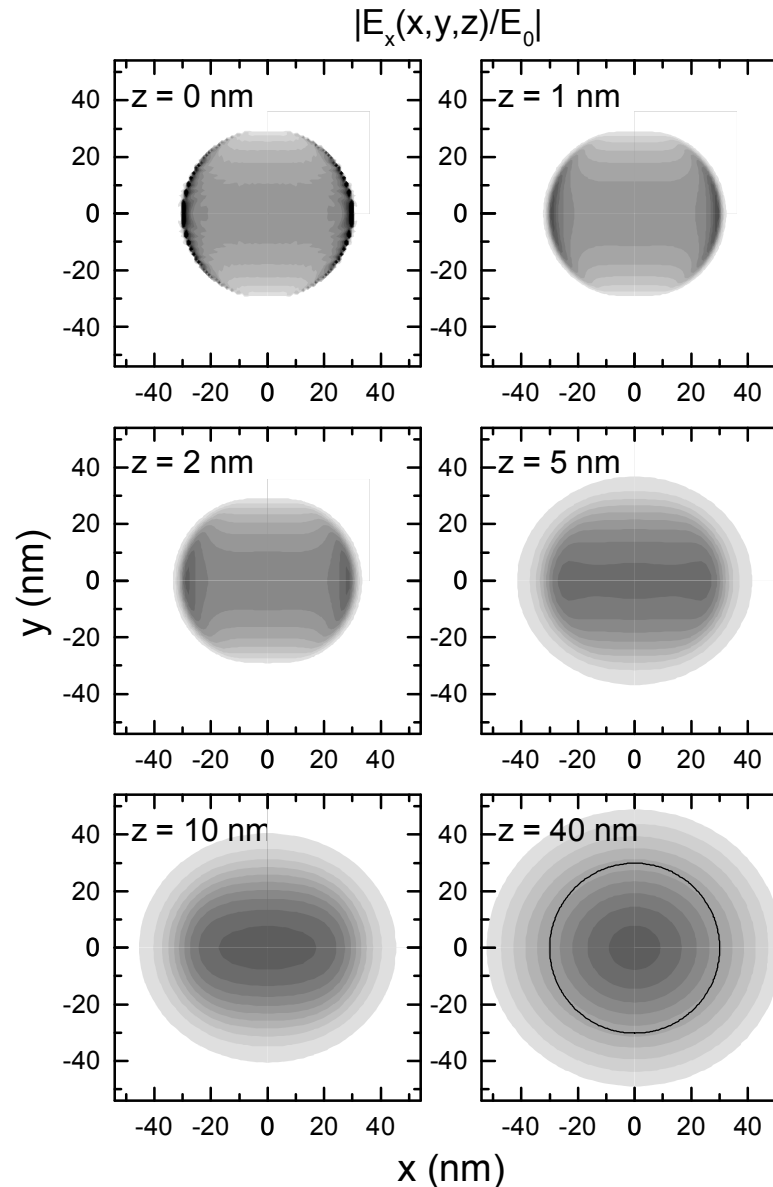
Monochromatic plane wave ($k_0 = 2\pi n_r / \lambda$ - $\lambda = 800$ nm) incident on slit (width $d = 50$ nm)



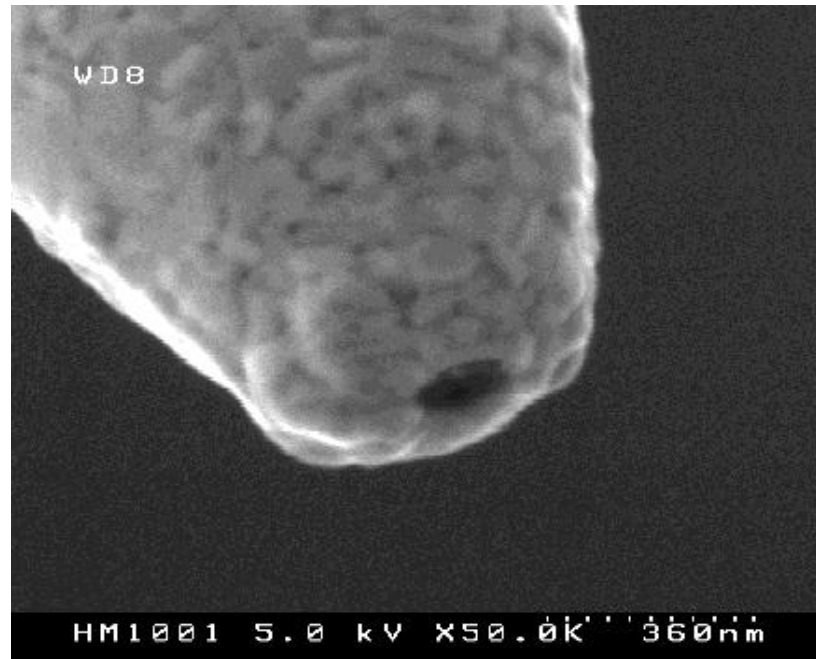
Evanescent waves: $k_{lat} > k_0 = 2\pi n_r / \lambda$ - $k_z^2 = (k_0^2 - k_{lat}^2) < 0$ - k_z imaginary

Propagating waves: $k_{lat} < k_0 = 2\pi n_r / \lambda$ - $k_z^2 = (k_0^2 - k_{lat}^2) > 0$ - k_z real

Two-dimensional
field distribution $E_x(x,y)$
behind a $r=30$ nm aperture
in a perfect metal film



Aperture-based near-field microscopy



- + excellent rejection of propagating waves
- low transmission efficiency (10^{-4} - 10^{-3} for 100 nm apertures)

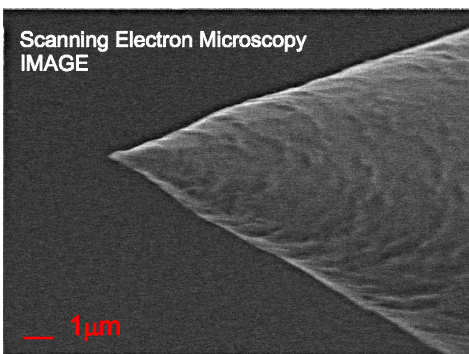
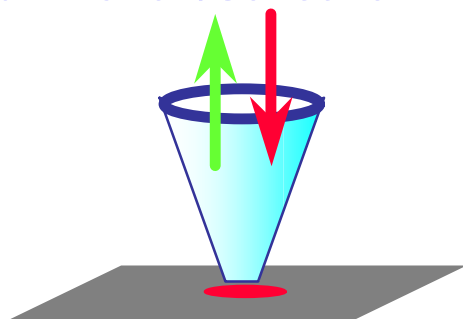
Introduction



Aperture-based near-field microscopy

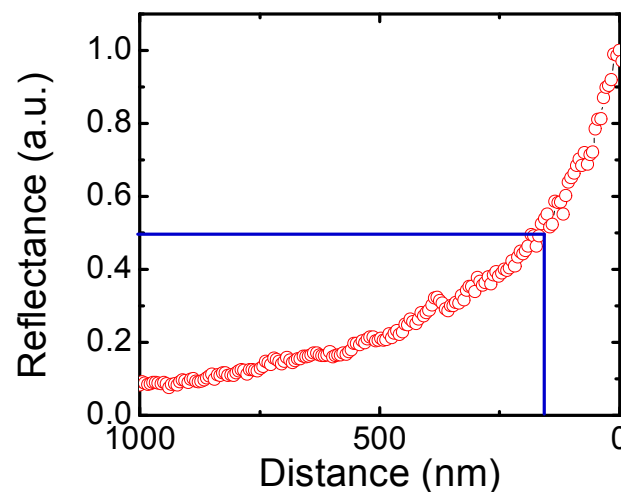
Near-field Reflectance Spectroscopy with Uncoated Fiber Probes

Illumination/Collection Mode



Transmission efficiency close to 1
Extremely high collection efficiency

Sensitive probing of near-field reflectance



Resolution down to $\lambda/5$ (150 nm)

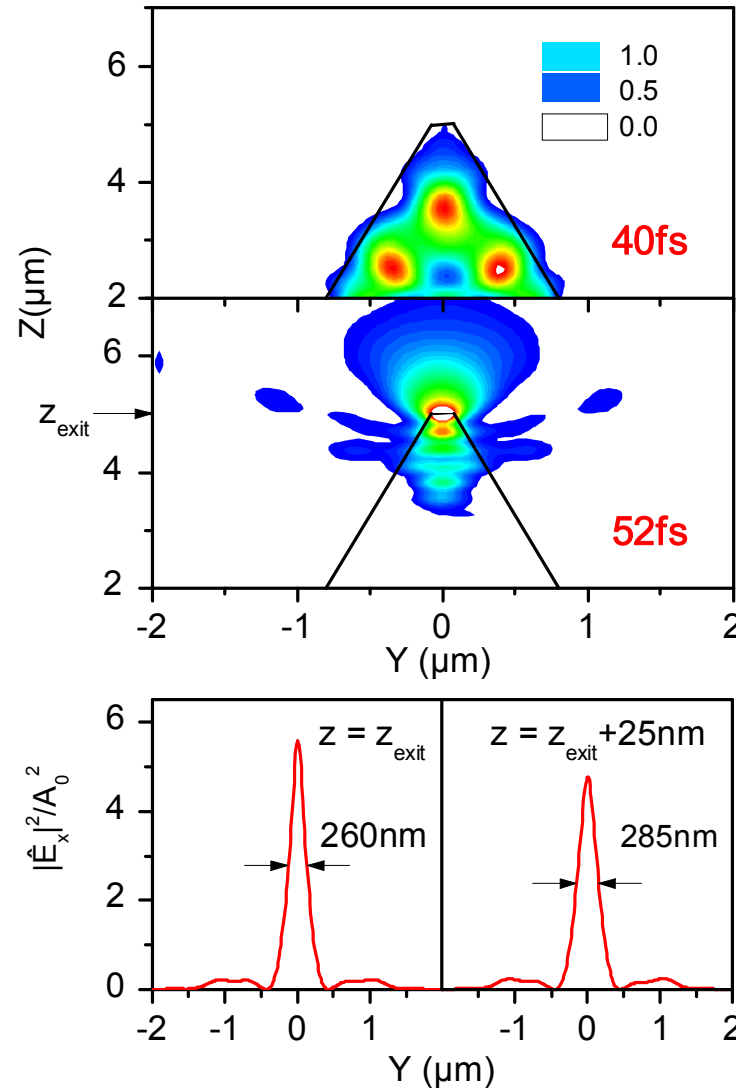
Experiment: F. Intonti et al, PRL **87**, 076801 (2001), PRB **63**, 075313 (2001)

Theory: R. Müller and C. Lienau, Appl. Phys. Lett. **76**, 3367 (2000).

Pulse propagation through uncoated fiber probes

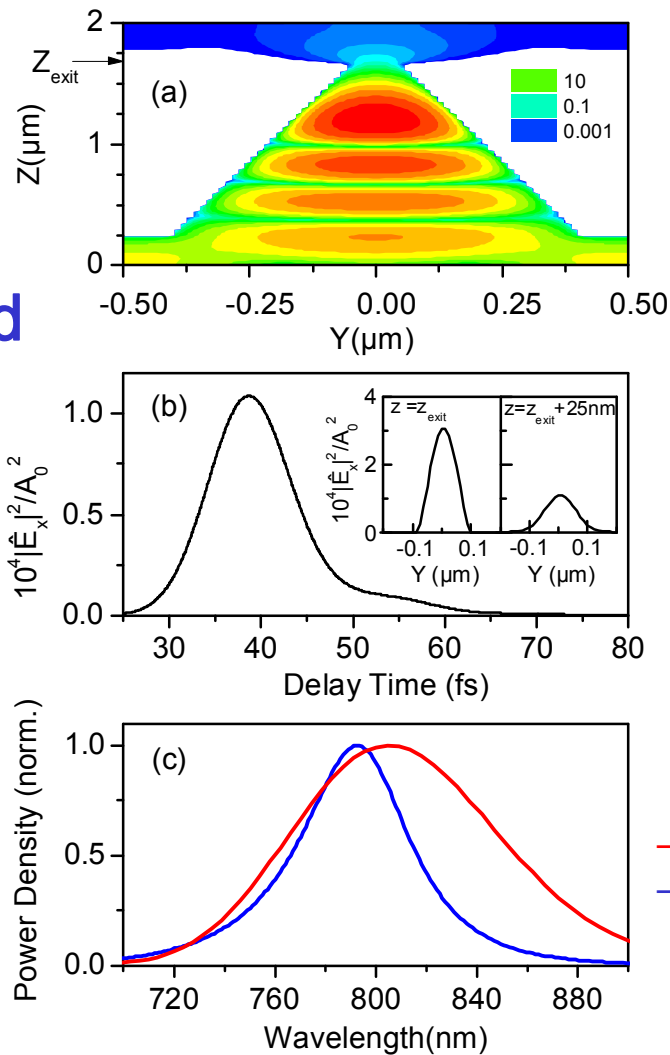
2D modelling
color-code
represents
 $|E|^2 A_0^2$

R. Müller Appl. Phys. Lett. 76, 3367 (2000)
and J. Microscopy 202, 339 (2001).



Pulse propagation through metal-coated fiber probes

R. Müller and C. Lienau
Appl. Phys. Lett.
76, 3367 (2000)
J. Microscopy 202,
339 (2001).



2D modelling
color-code represents
 $|E|^2 A_0^2$

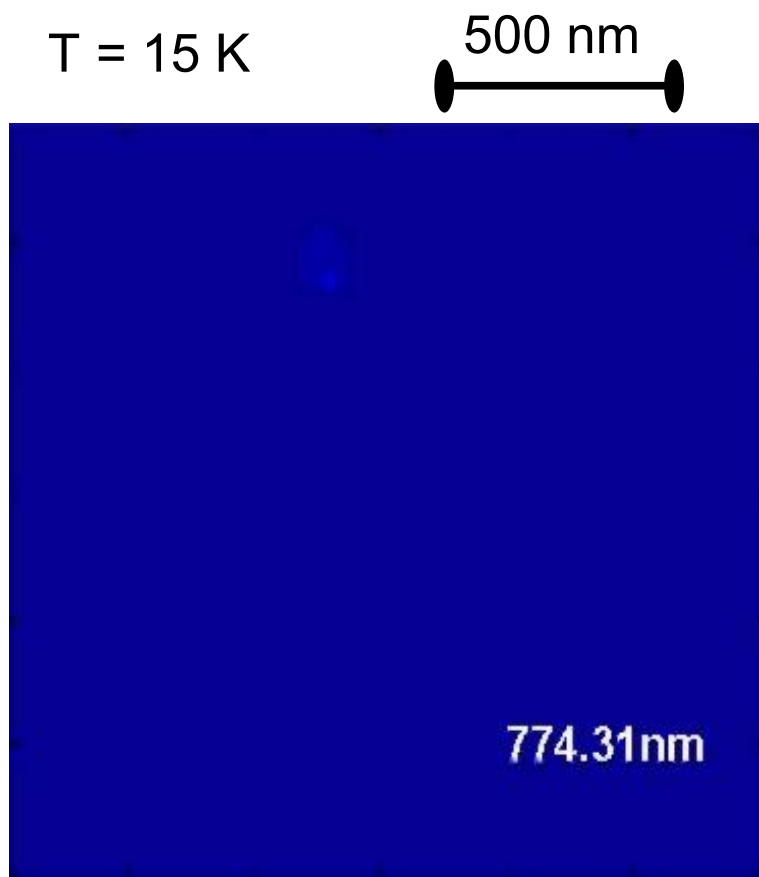
$|E|^2 A_0^2$ of a
10 fs Gaussian pulse
at the 100 nm aperture

— input pulse
— output pulse

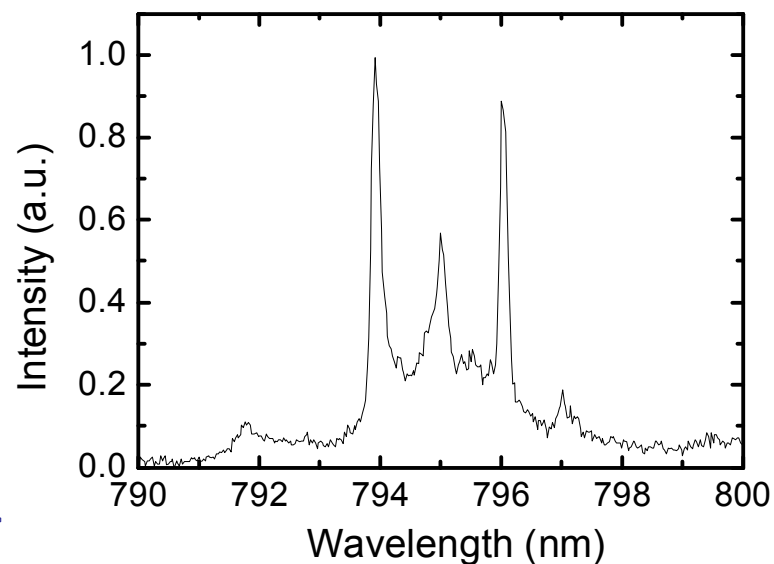
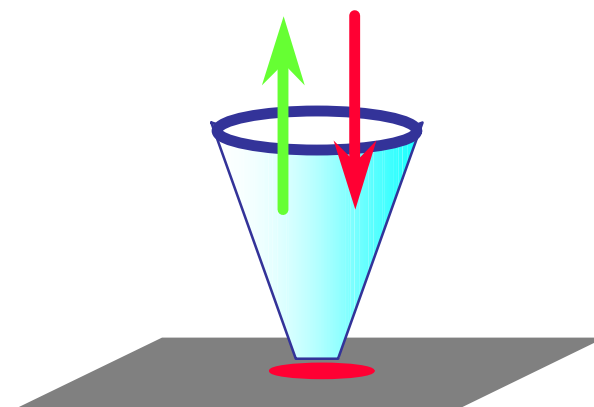
Application: Raman spectroscopy of Carbon Nanotubes

A. Hartschuh et al., Phys. Rev. Lett. (2003)

Nanophotoluminescence of single localized excitons



F. Intonti et al., Phys. Rev. Lett. 87, 076801 (2001).



Apertureless Near-field Scanning Optical Microscopy

- + Strong field enhancement (10 x) at ultrasharp metal tips
- + Spatial resolution limited by radius of curvature of the tip
- + Spatial resolution down to 10 nm (and beyond?)

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RW laser

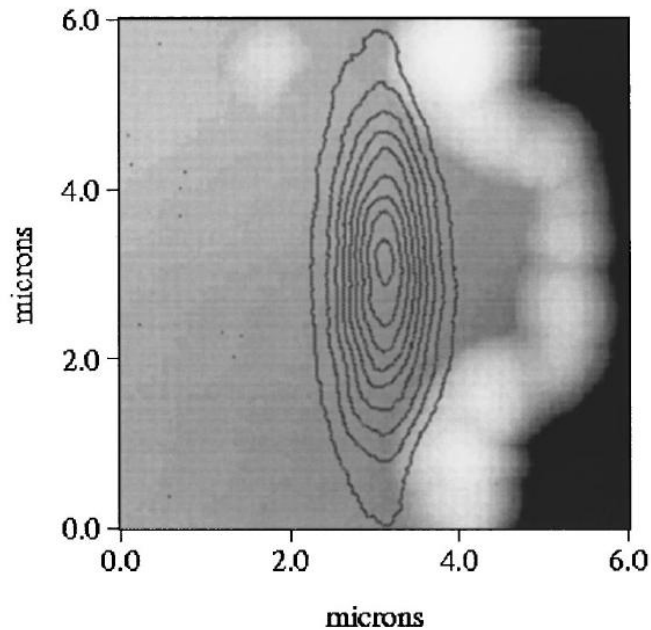


FIG. 1. Near-field collection mode image of the GRINSCH laser diode emission intensity (constant intensity contours) superimposed on the topographical shear-force image (greyscale). The metal contact on top of the mesa slightly protrudes over the facet of the device, yielding the brighter topographic regions.

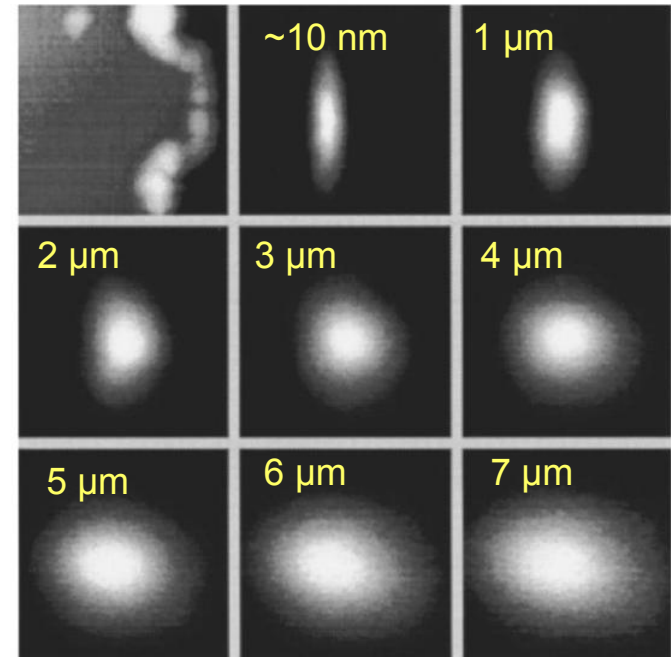
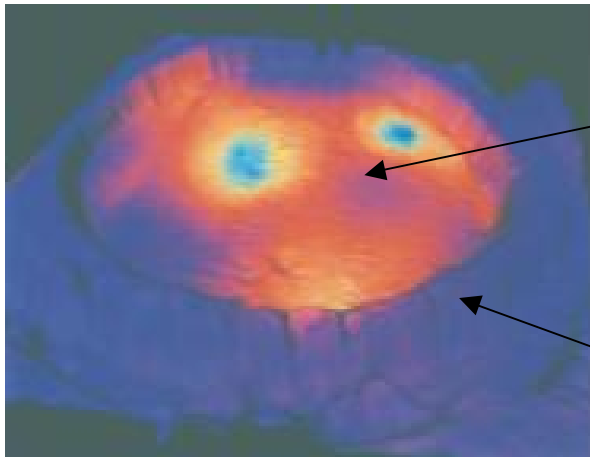


FIG. 3. Images of the beam propagation from the near field to $7\ \mu\text{m}$ from the laser diode taken in $1\ \mu\text{m}$ steps. The top left image is the topography to provide the physical orientation for the laser emission images given in left to right sequence. All images are $6 \times 6\ \mu\text{m}$. The greyscale for each image is chosen to maximize contrast.

W. D. Herzog, M. S. Ünlü, B. B. Goldberg, G. H. Rhodes, and C. Harder, *Appl. Phys. Lett.* 65 688 (1997).

VCSEL

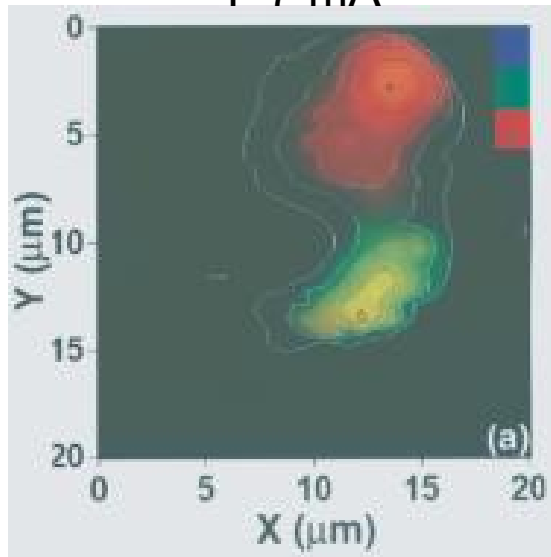


Emission intensity (color)

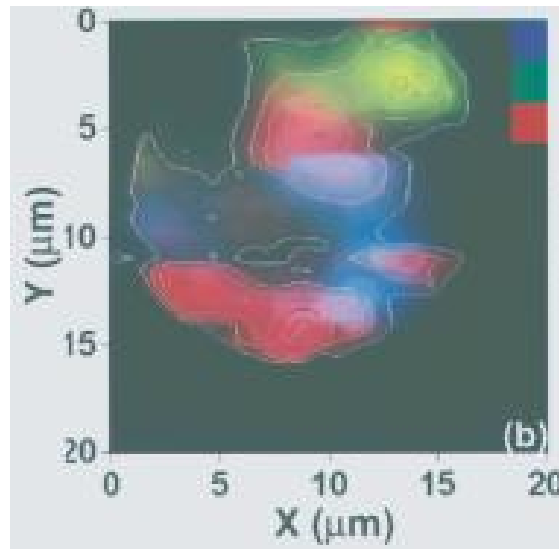
Van der Rhodes et al.
Appl. Phys. Lett. 72,1811 (1998)

VCSEL aperture (blue)

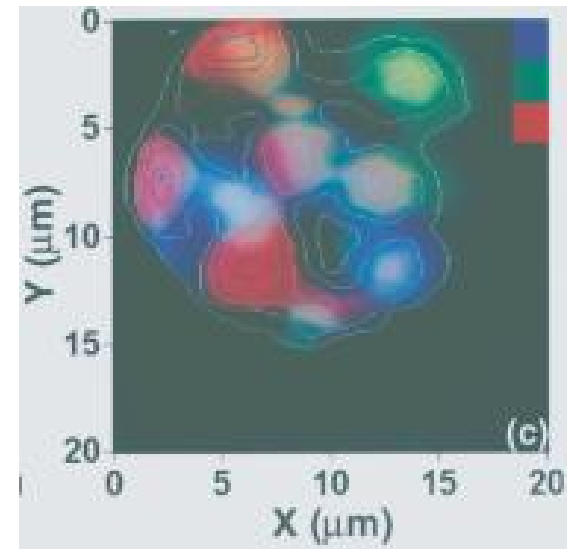
$I=7$ mA



$I=10$ mA



$I=15$ mA



Such work at devices is extremely useful,

1. if confocal microscopy fails

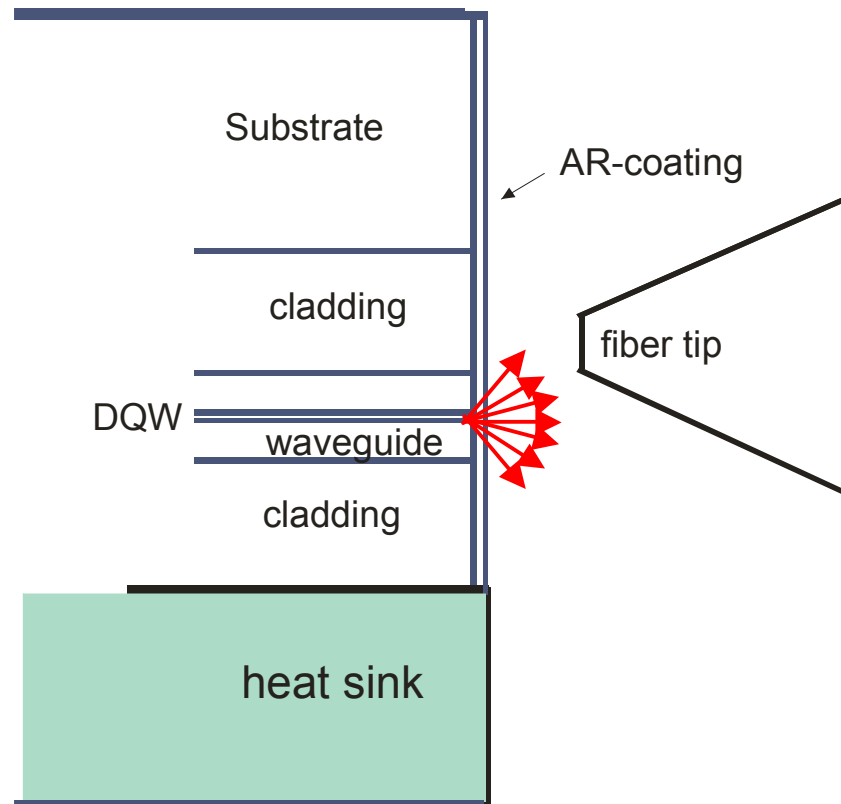
if not, you waste your resources

2. if the fiber tip does not influence the emission

unfortunately, this is the case when investigating lasing devices

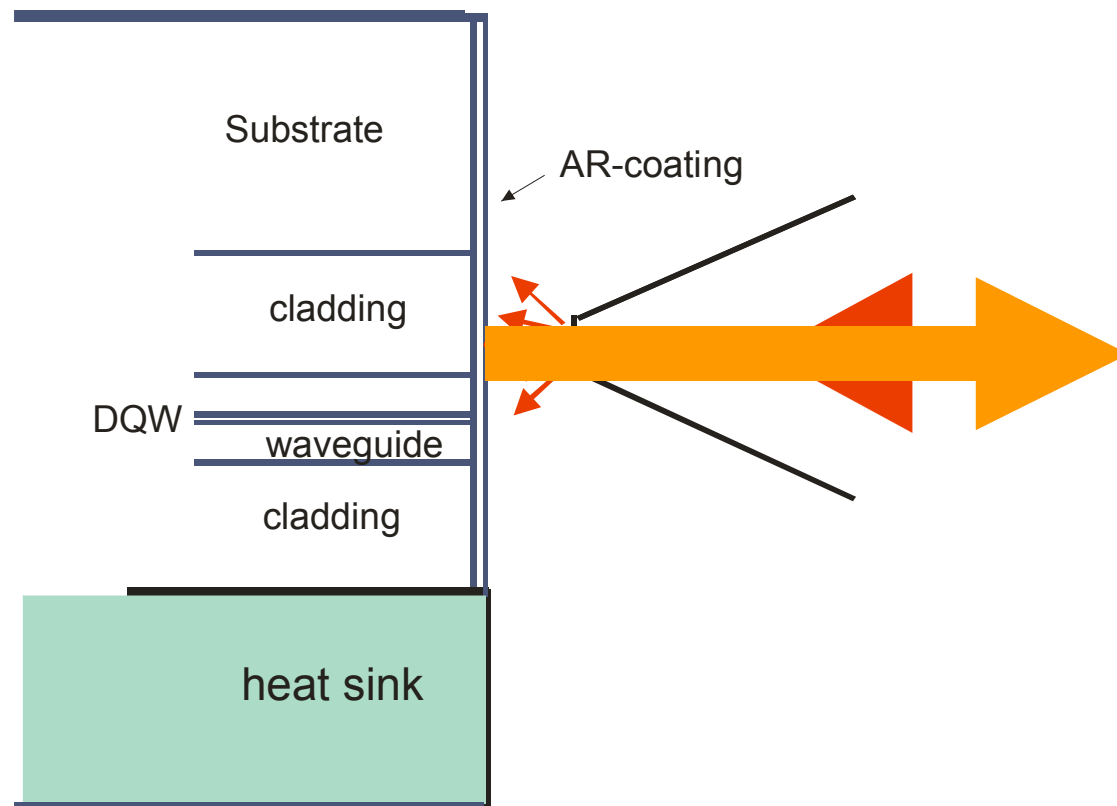
These statements hold in general for the application of NSOM !

- Spontaneous emission



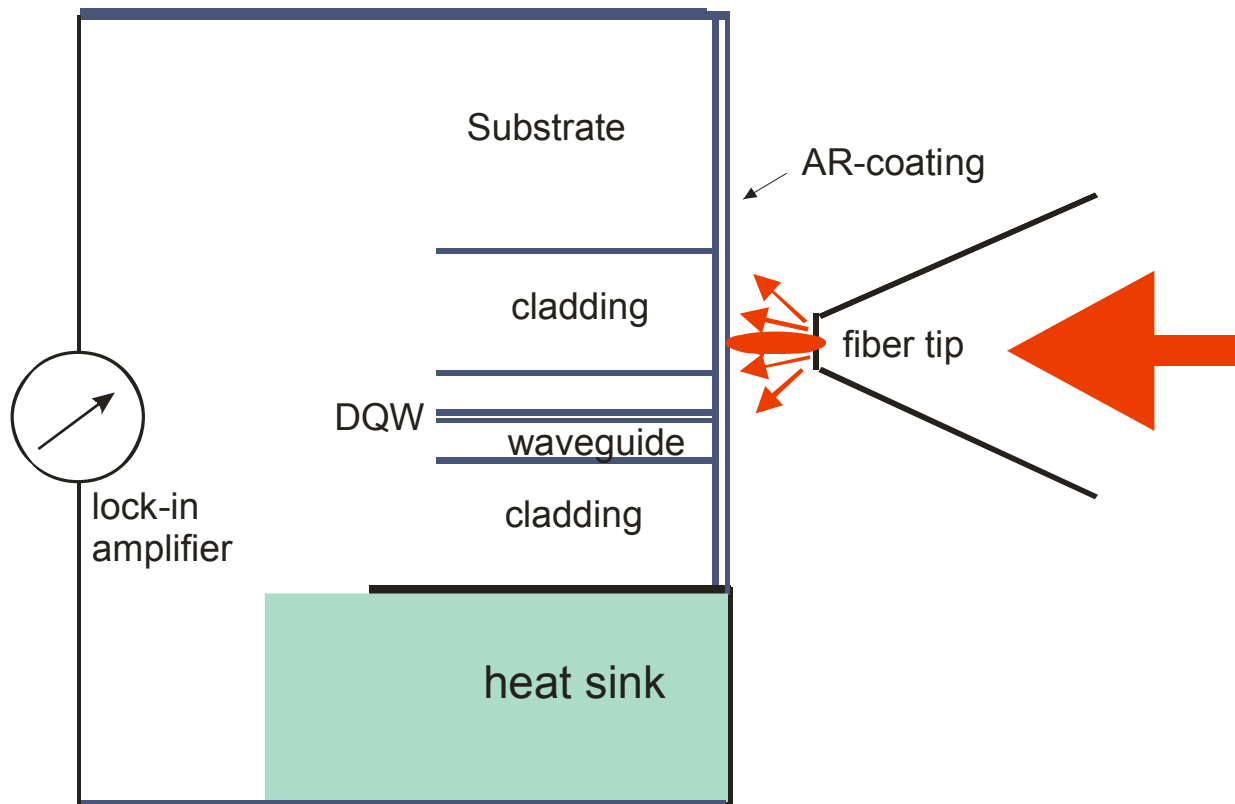
What is better?

- Spontaneous emission
- Photoluminescence



What is better?

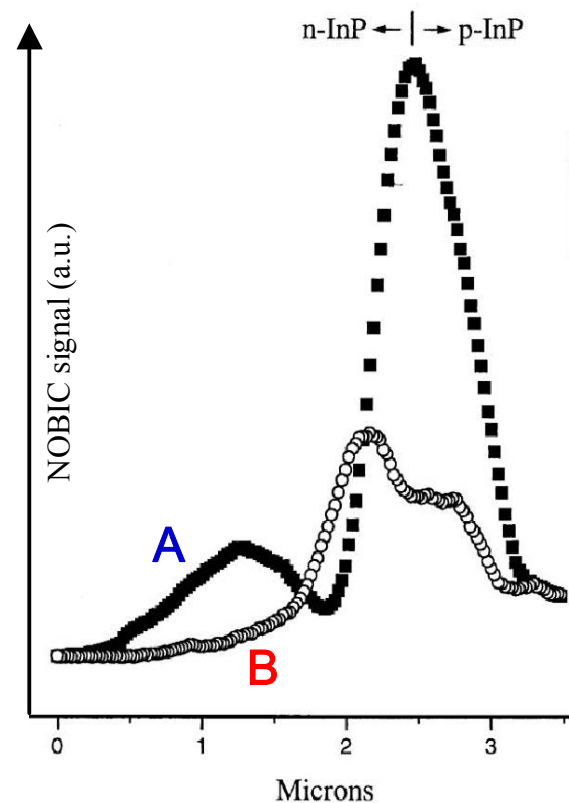
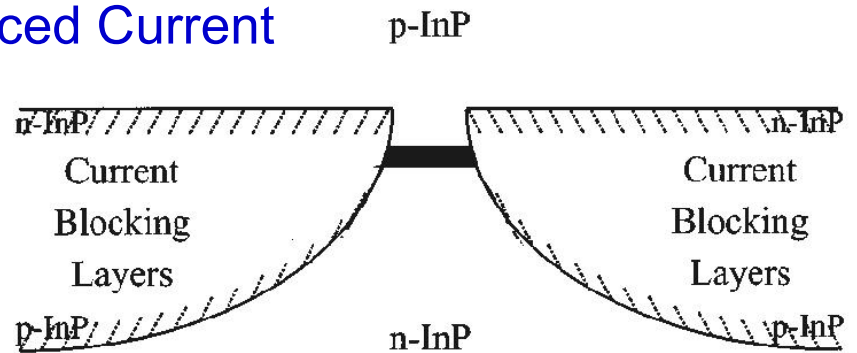
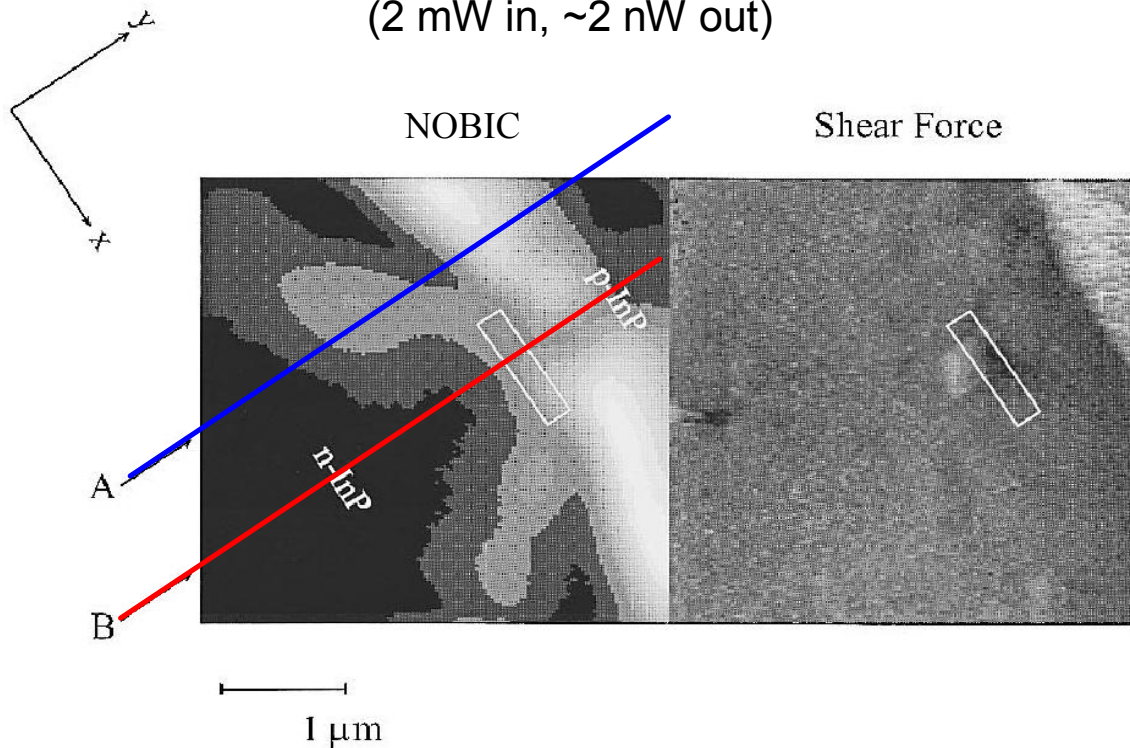
- Spontaneous emission
- Photoluminescence
- Absorption (Photocurrent) NOBIC



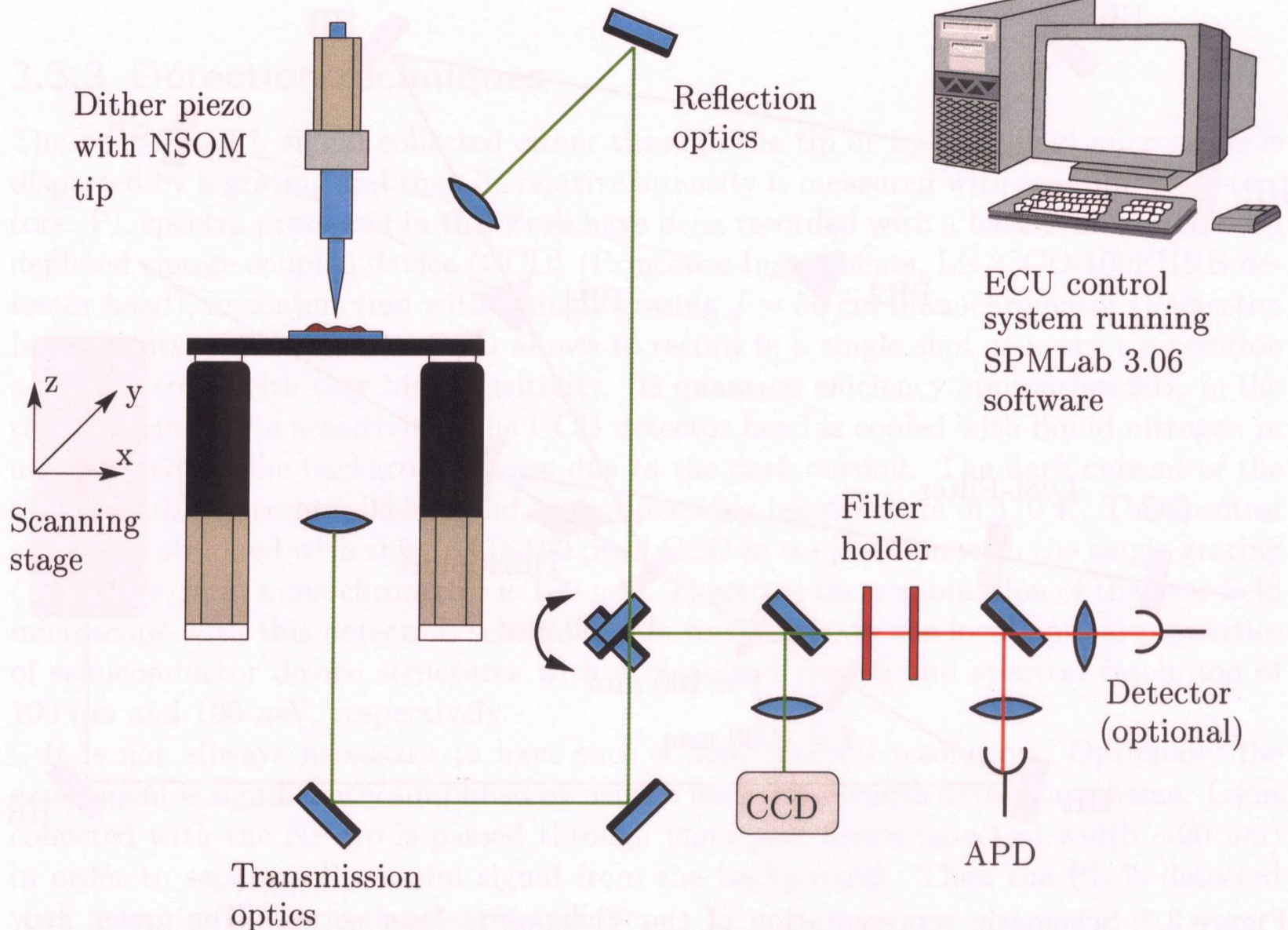
NOBIC = Nearfield Optical Beam Induced Current

Buratto et al. (AT&T Bell Lab.)
Appl. Phys. Lett. 65, 2654 (1994)

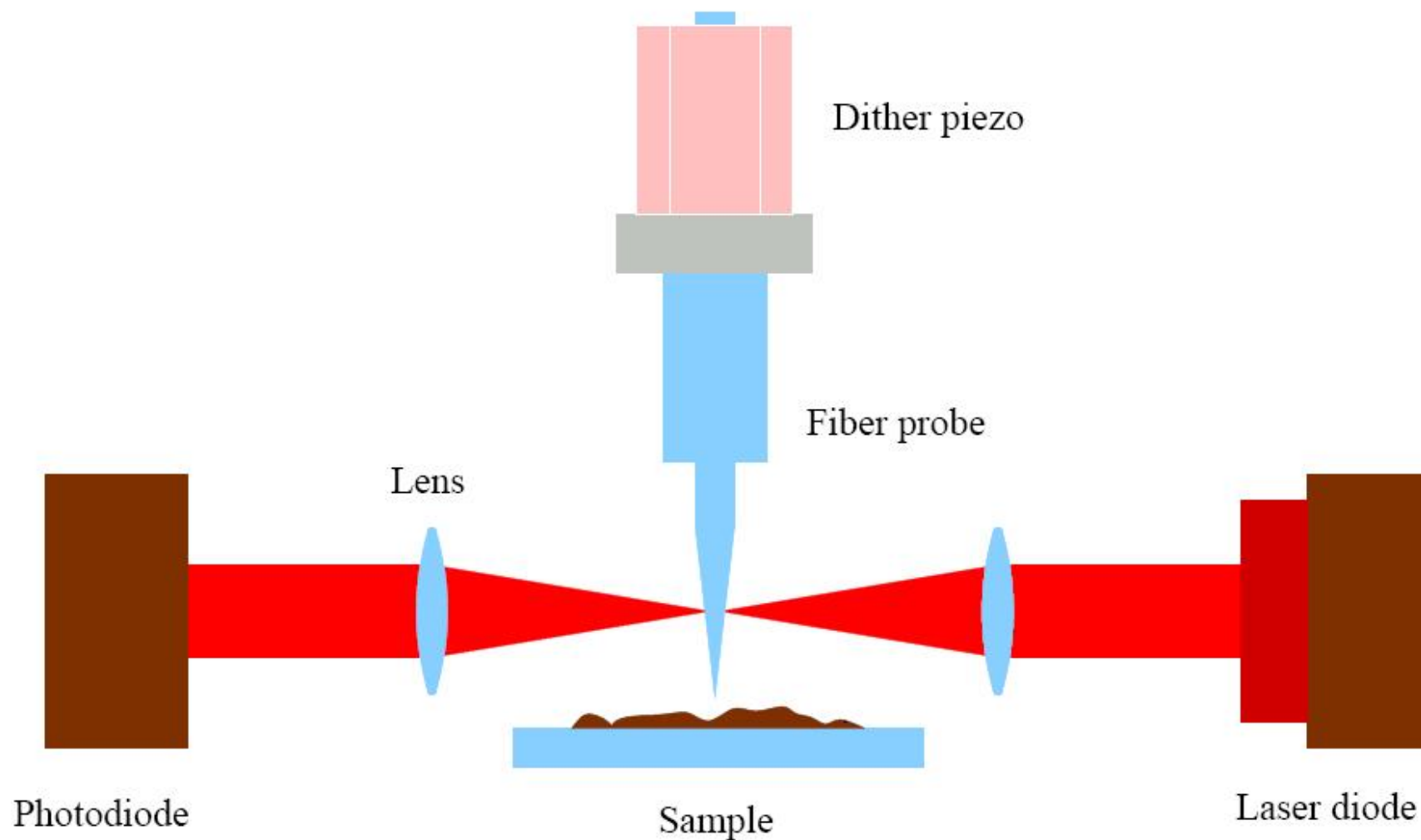
Excitation at 633 nm
(2 mW in, ~2 nW out)



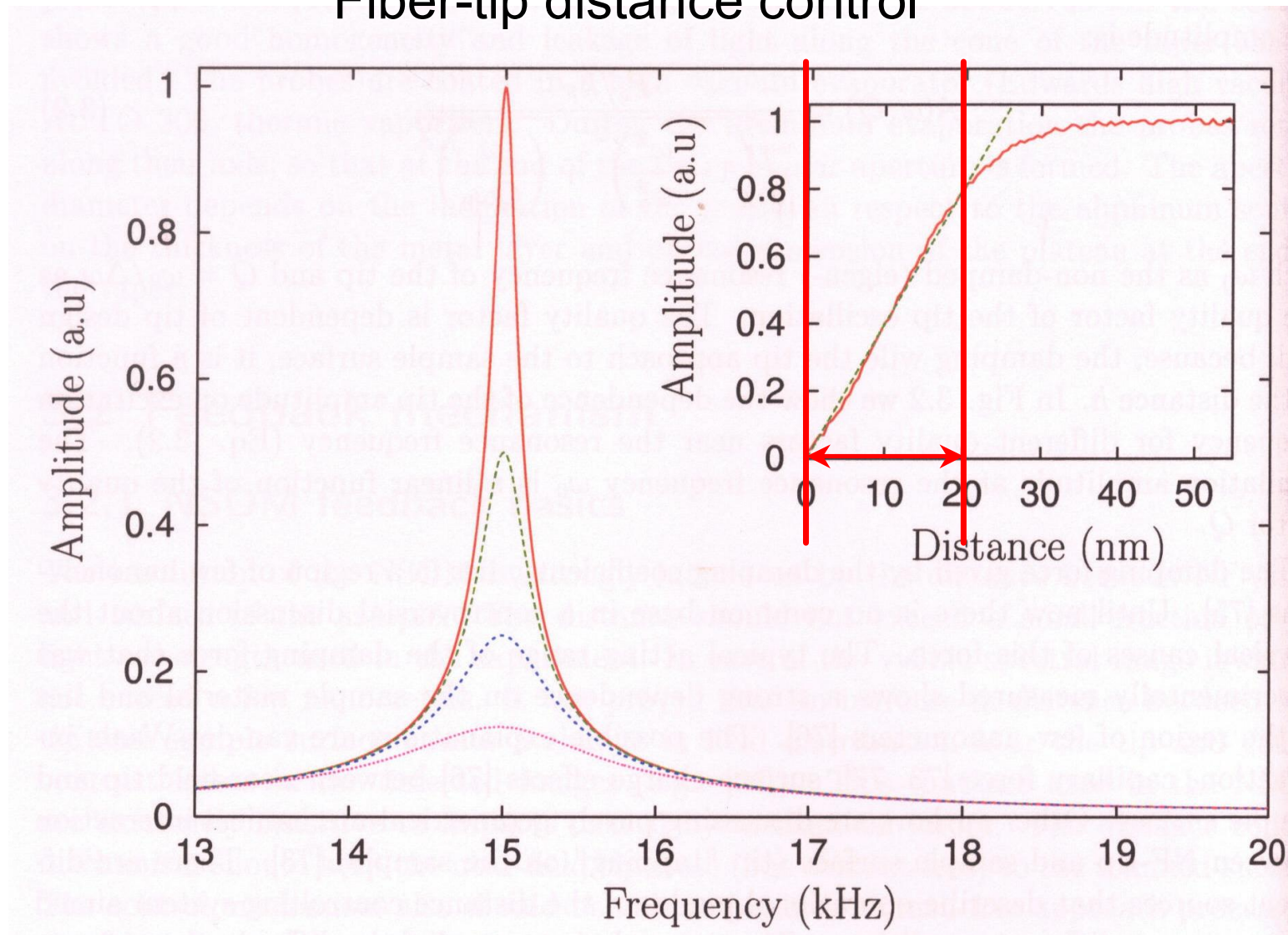
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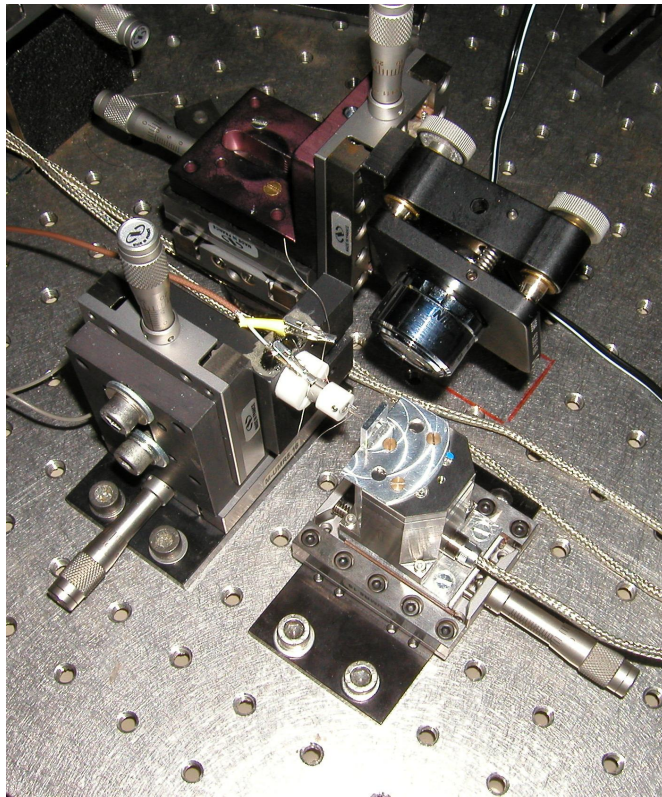
Fiber-tip distance control



Fiber-tip distance control

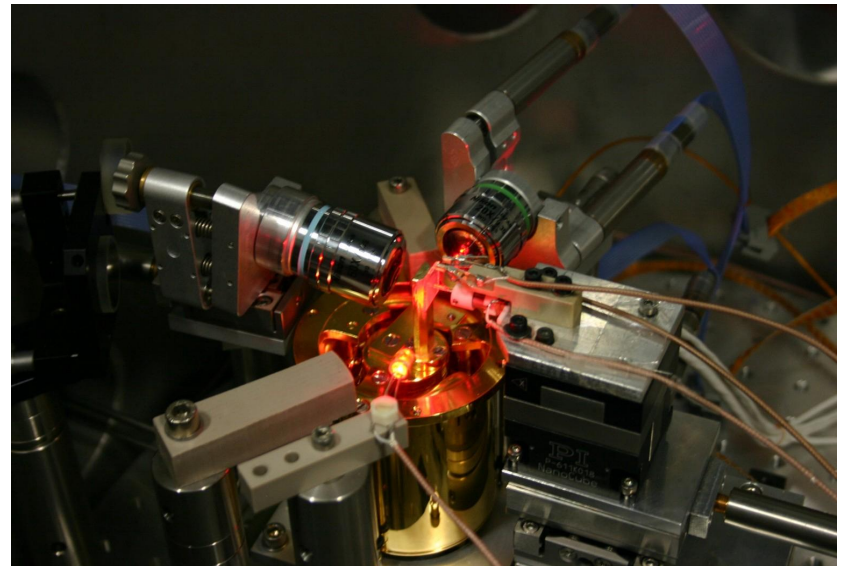


Home-built Scanning Near-field Optical Microscopes



G. Behme et al., Rev. Sci. Instrum.
68, 3458 (1997)

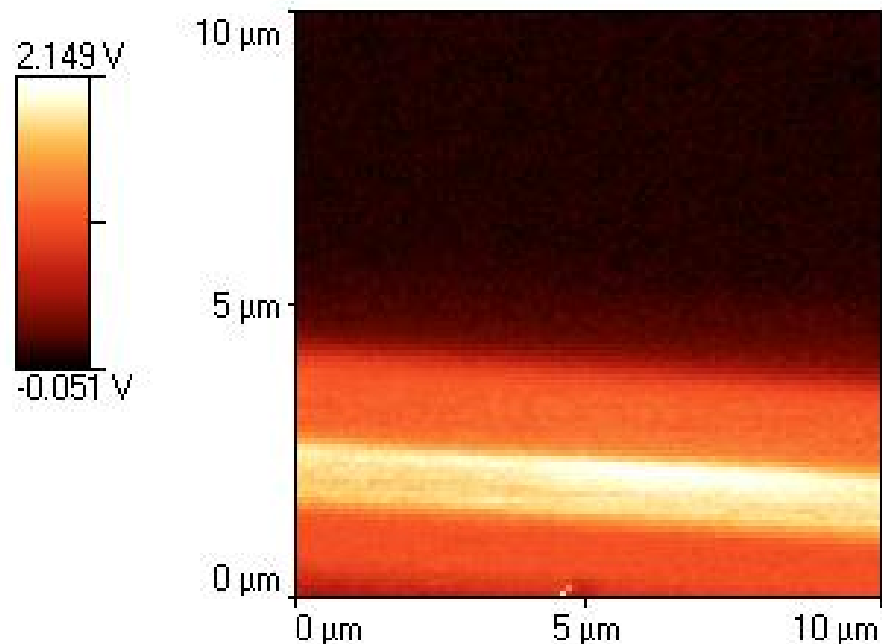
Set-up with cryostat for
Measurements at 10 - 300 Kelvin



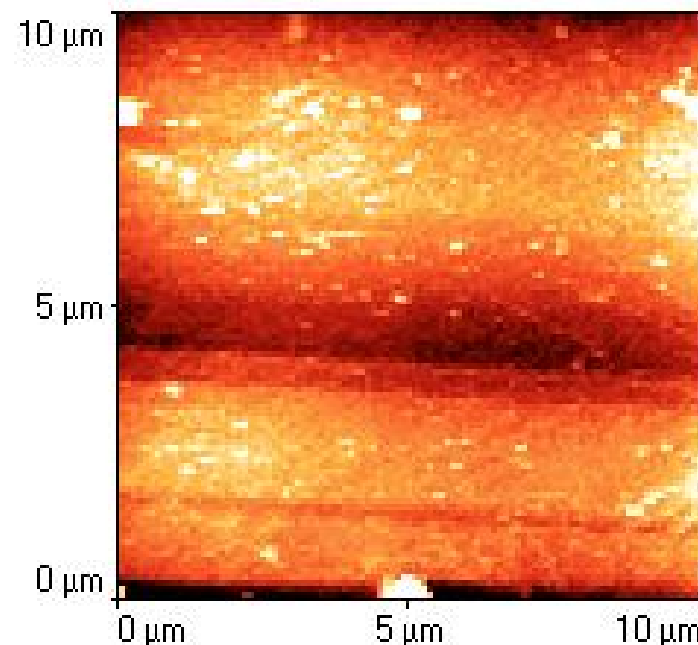
"Customers:"
Uni Magdeburg, D
KTH Stockholm, Sweden
Forschungszentrum Jülich, D
University of Arkansas, USA

NSOM data

PC-signal (R)

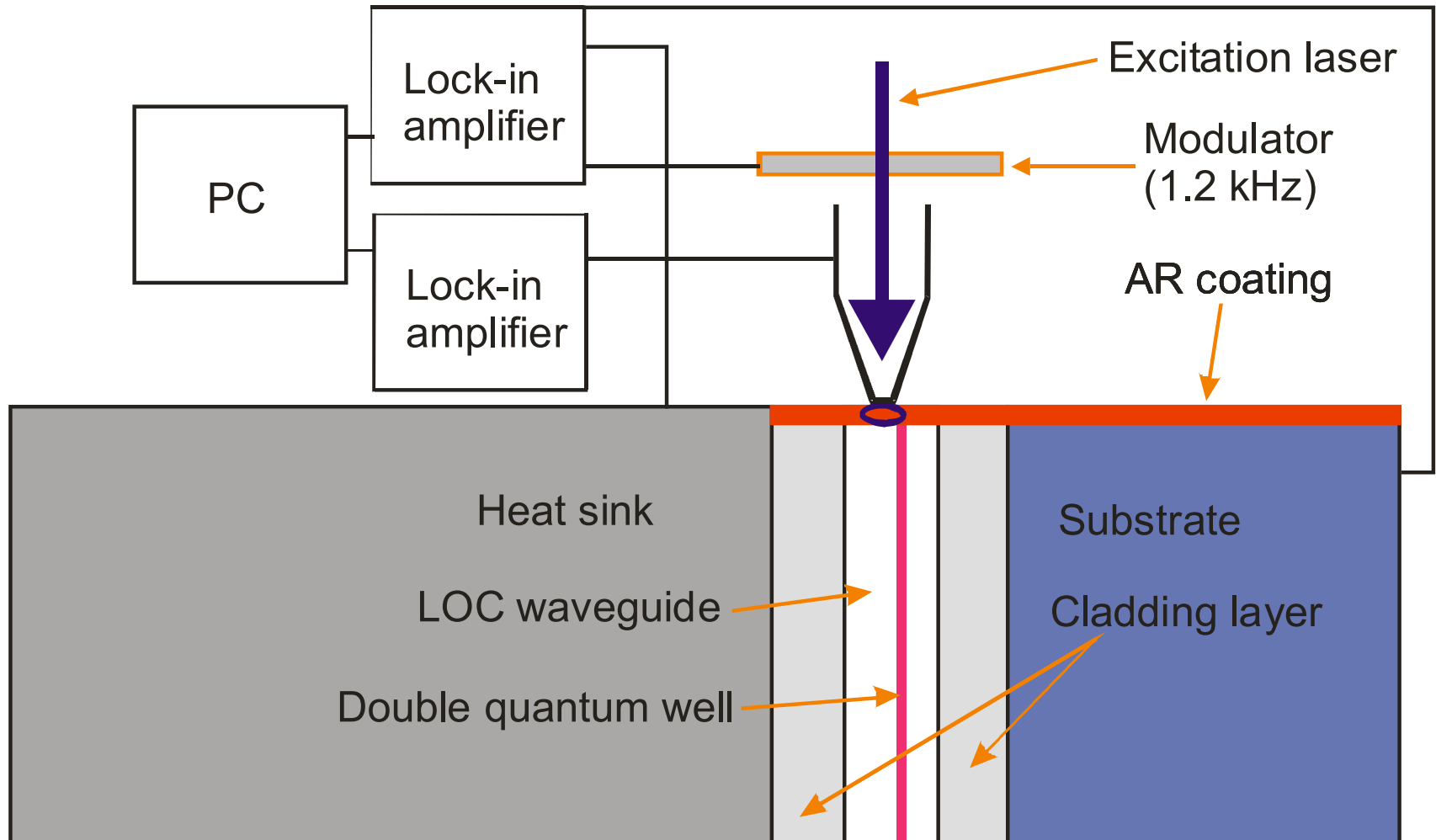


topography



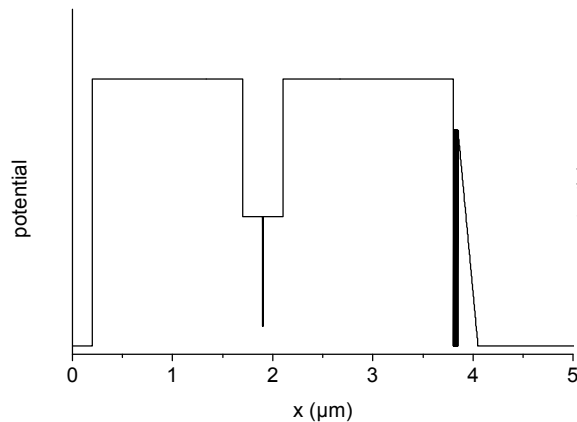
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Setup

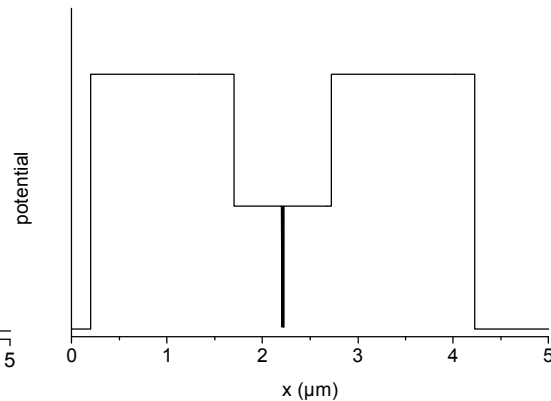


Results obtained at single symmetric and asymmetric waveguides

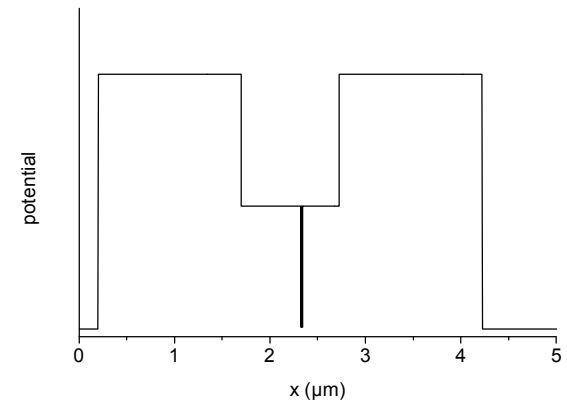
Step INDEX
SIN



Large Optical Cavity
LOC **symmetric** waveguide

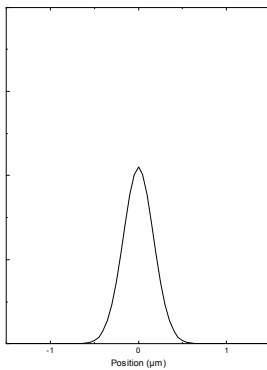
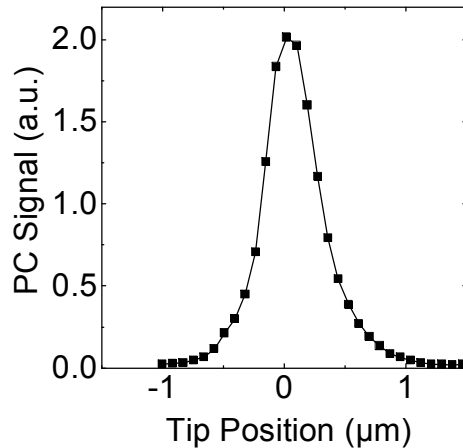


Large Optical Cavity
LOC **asymmetric** waveguide

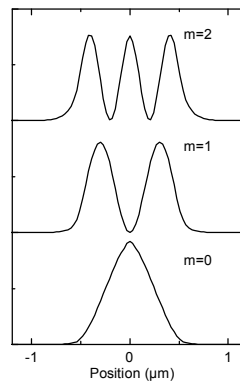
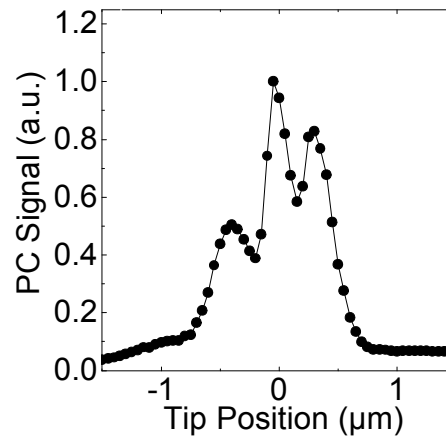


resonant NSOM Photocurrent (Absorption)

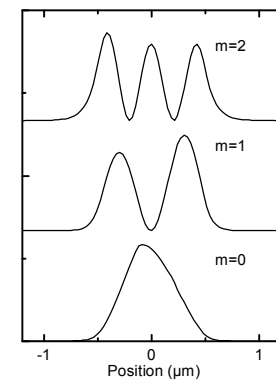
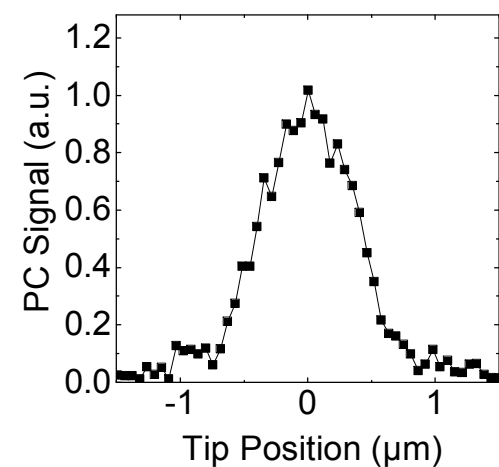
Step INDEX SIN



Large Optical Cavity LOC **symmetric** waveguide



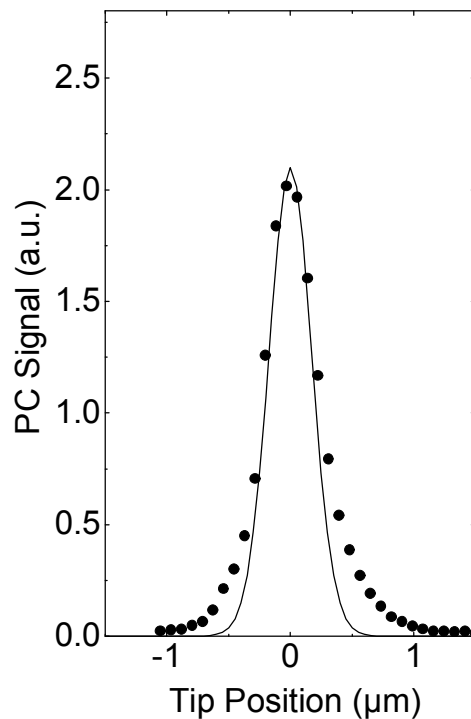
Large Optical Cavity LOC **asymmetric** waveguide



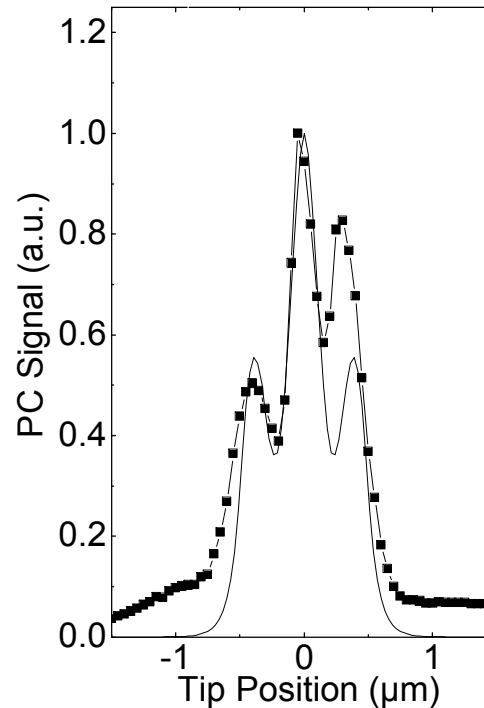
Comparison of the guided modes in a waveguide and the wavefunctions in a quantum well

Field	$\Psi(\mathbf{r}, t) = \Psi(\mathbf{r}) e^{i\omega t}$ Wave function in Bloch form	$\mathbf{H}(\mathbf{r}, t) = \mathbf{H}(\mathbf{r}) e^{i\omega t}$ Harmonic modes in Bloch form
Eigenvalue problem	$H\Psi = E\Psi$	$\Theta\mathbf{H} = (\omega/c)^2\mathbf{H}$
Hermitian operator	$H = \frac{-\hbar^2\nabla^2}{2m} + V(\mathbf{r})$	$\Theta = \nabla \times \left(\frac{1}{\varepsilon(\mathbf{r})} \nabla \times \right)$
Periodicity of the system	$V(\mathbf{r}) = V(\mathbf{r} + \mathbf{R})$ Periodic potential for all lattice vectors $\mathbf{R}$	$\varepsilon(\mathbf{r}) = \varepsilon(\mathbf{r} + \mathbf{R})$ Periodic dielectric for all lattice vectors $\mathbf{R}$
Variational theorem	$E_{var} = \frac{\langle \Psi H \Psi \rangle}{\langle \Psi \Psi \rangle}$ is minimized when Ψ is an eigenstate of H	$E_{var} = \frac{(\mathbf{H}, \Theta\mathbf{H})}{(\mathbf{H}, \mathbf{H})}$ is minimized when $\mathbf{H}$ is normal mode of Θ

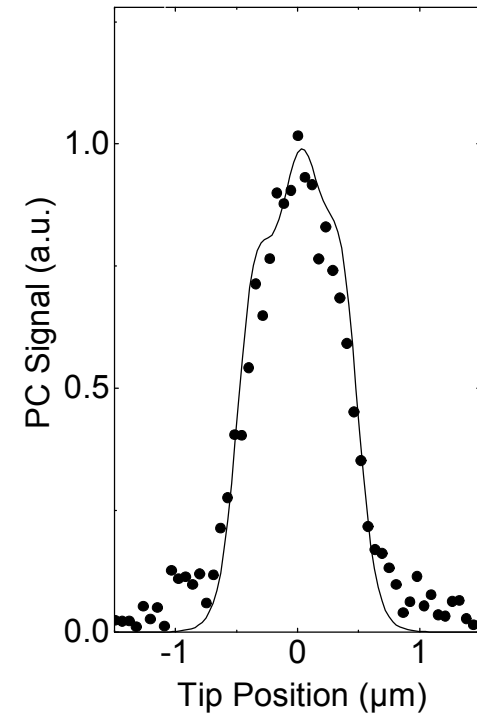
Step Index
SIN



Large Optical Cavity
LOC **symmetric** waveguide



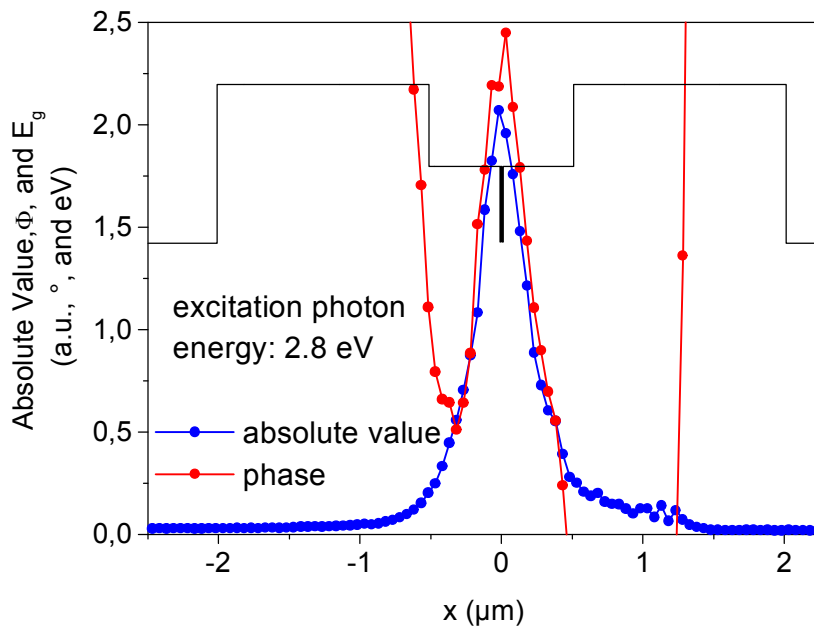
Large Optical Cavity
LOC **asymmetric** waveguide



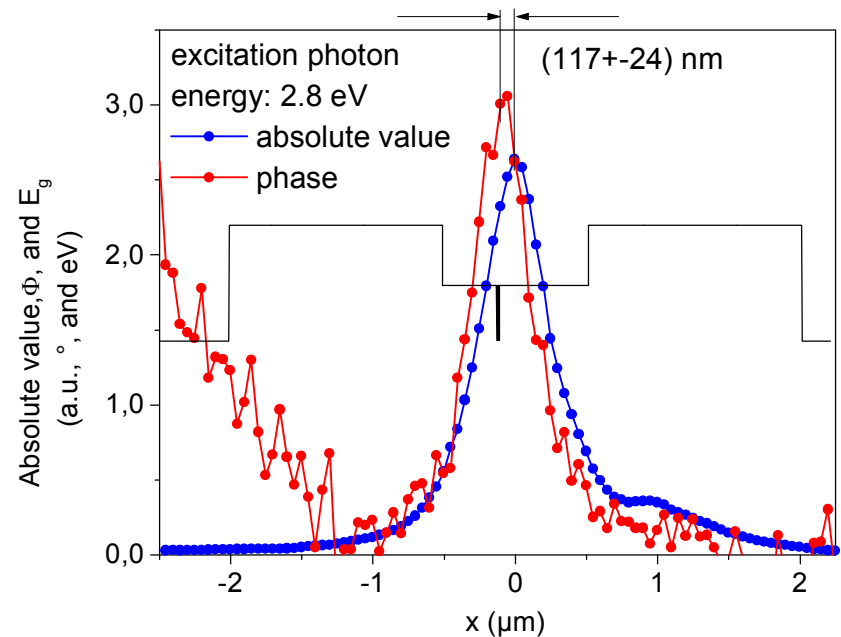
- NSOM Photocurrent reveals mode structure of waveguides
- asymmetric waveguides have a specific NSOM Photocurrent signature

NOBIC contrast from the photocurrent phase for non-resonant (surface) excitation

NOBIC line scans from the symmetric LOC structure



NOBIC line scans from the asymmetric LOC structure



Single waveguides

- Resonant excitation (coupling into the waveguide modes)
- NOBIC reveals mode structure of waveguides
- asymmetric waveguides have a specific NOBIC signature
- Non-resonant excitation (pure carrier-pair creation)
- detection of the QW within the waveguide

T. Guenther et al.

„Near-Field Photocurrent Imaging of the Optical Mode Profiles of Semiconductor Laser Diodes“

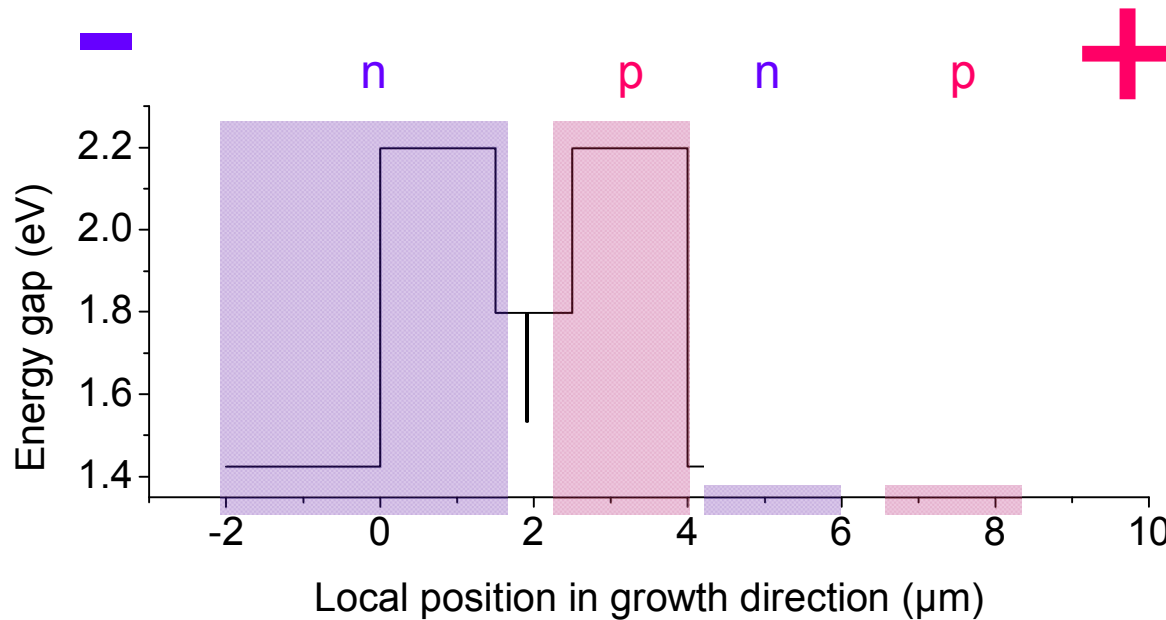
Appl. Phys. Lett. **78** 1463-1465 (2001).

Nanostacks[®] - interband cascade diode laser

'regular' high-power diode laser array



Nanostack[®]



J. P. van der Ziel and W. T. Tsang, *Appl. Phys. Lett.* **41**, 499-501 (1982).

J. Ch. Garcia, E. Rosencher, Ph. Collot, N. Laurent, J. L. Guyaux, B. Vinter, and J. Nagle, *Appl. Phys. Lett.* **71**, 3752-3754 (1997).

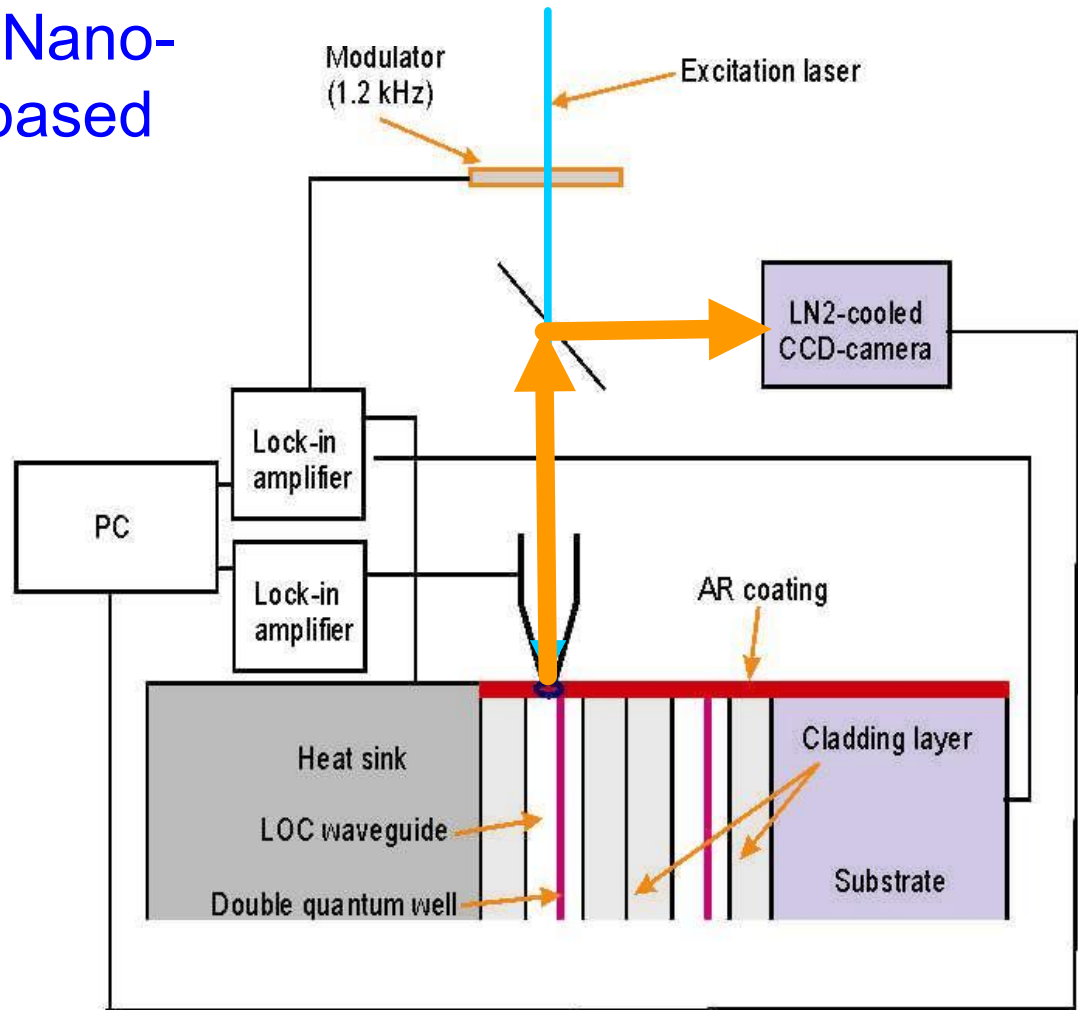
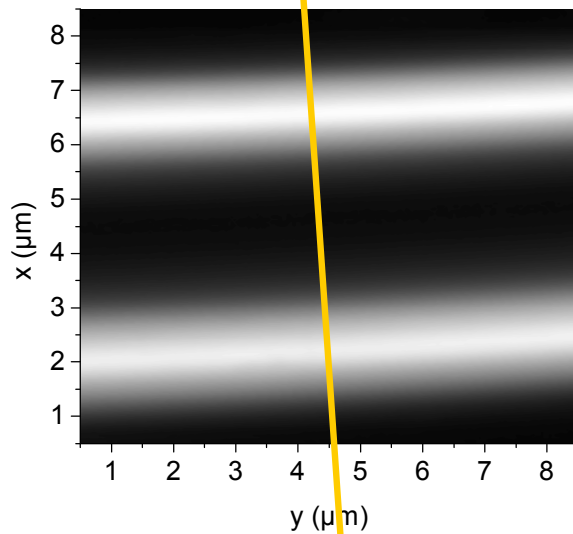
S. G. Patterson, G. S. Petrich, R. J. Ram, and L. A. Kolodziejski, *Electron Lett.* **35**, 395 (1999).

C. Hanke, L. Korte, B. D. Acklin, M. Behringer, G. Herrmann, J. Luft, B. De Odorico, M. Marchiano, J. Wilhelmi, *Proc. SPIE* **3947**, 50-57(2000).

http://www.osram-os.com/news/news_power.html

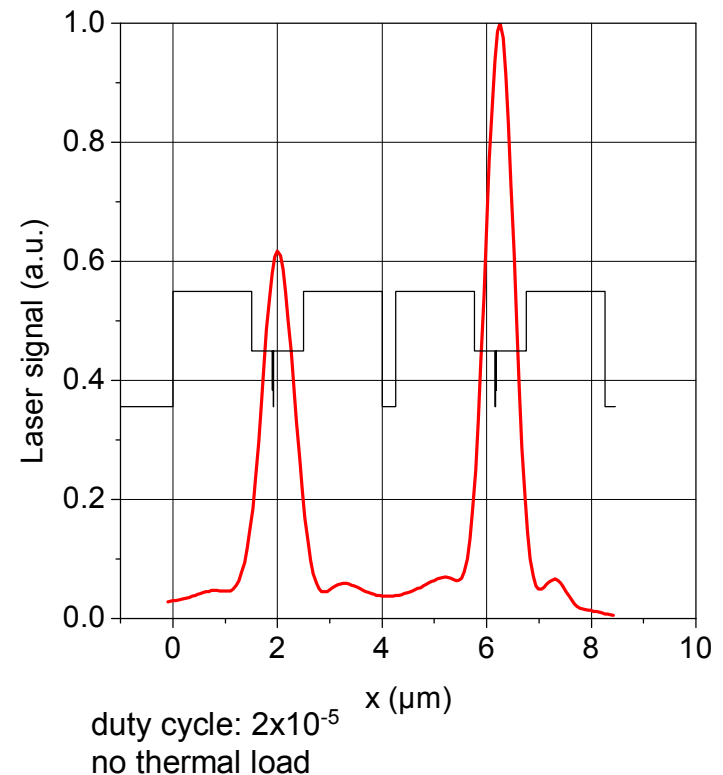
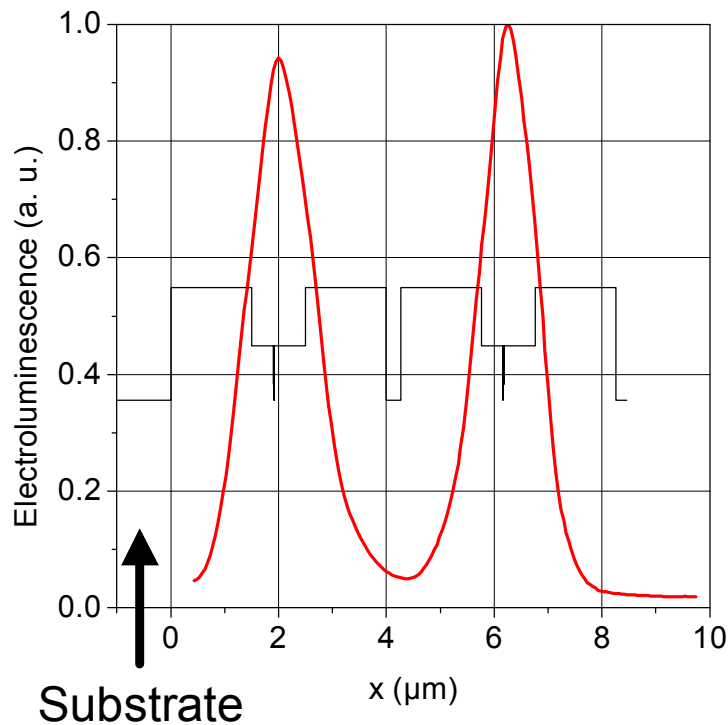
Detailed analysis of Nano-stacks® by NSOM-based spectroscopy

- Photoluminescence
- Photocurrent (Absorption)
- Device emission

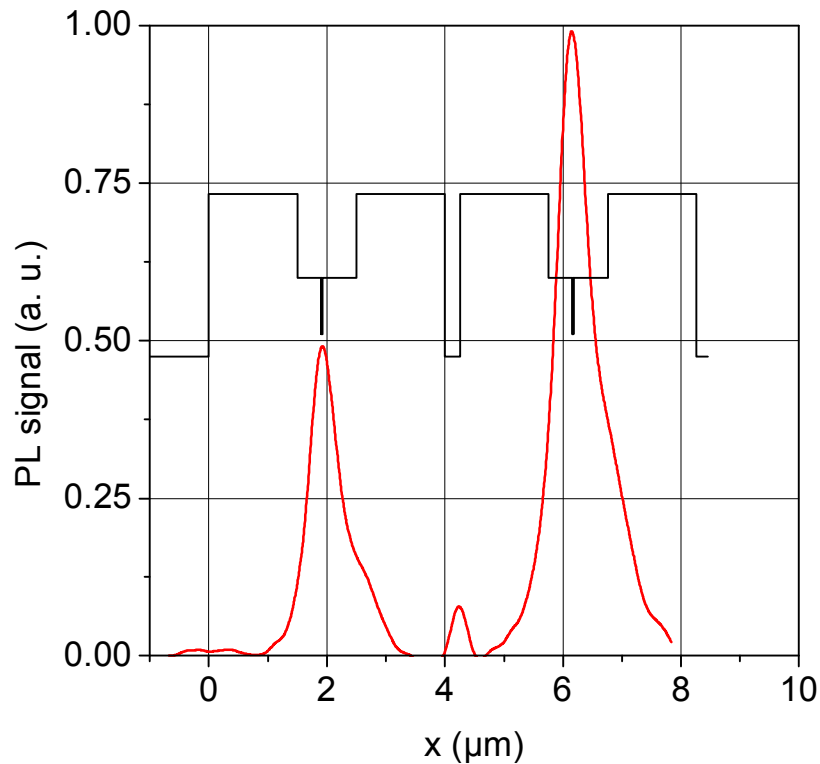


Results from Nanostacks® with two vertically stacked asymmetric waveguides

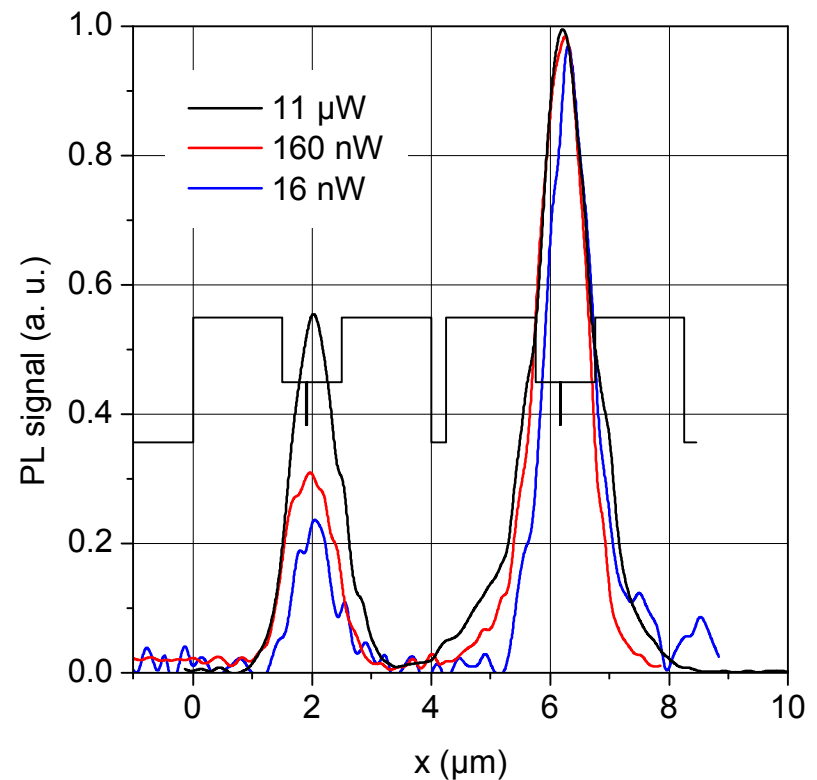
NSOM Device Emission



NSOM PL Emission



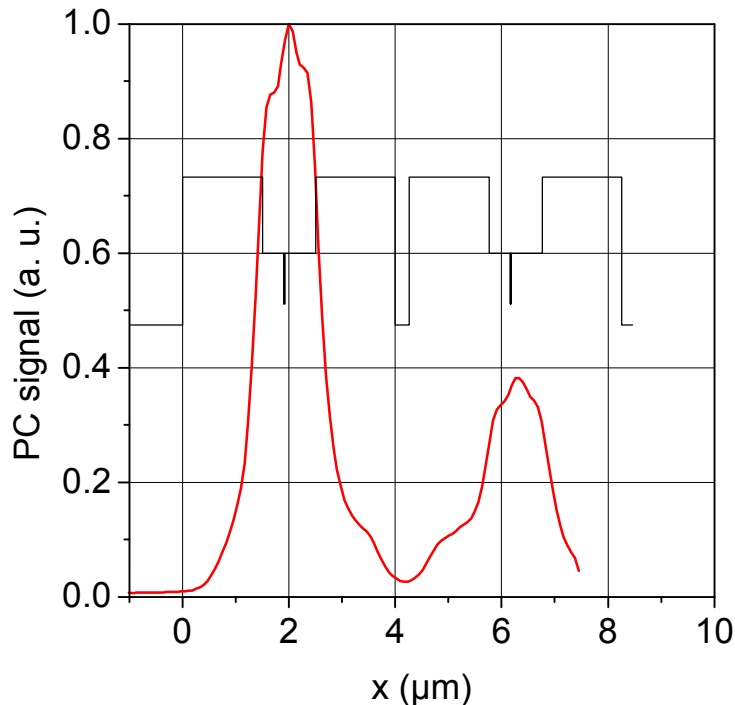
excitation photon energy: 2.8 eV
detection photon energy: 1.5 eV



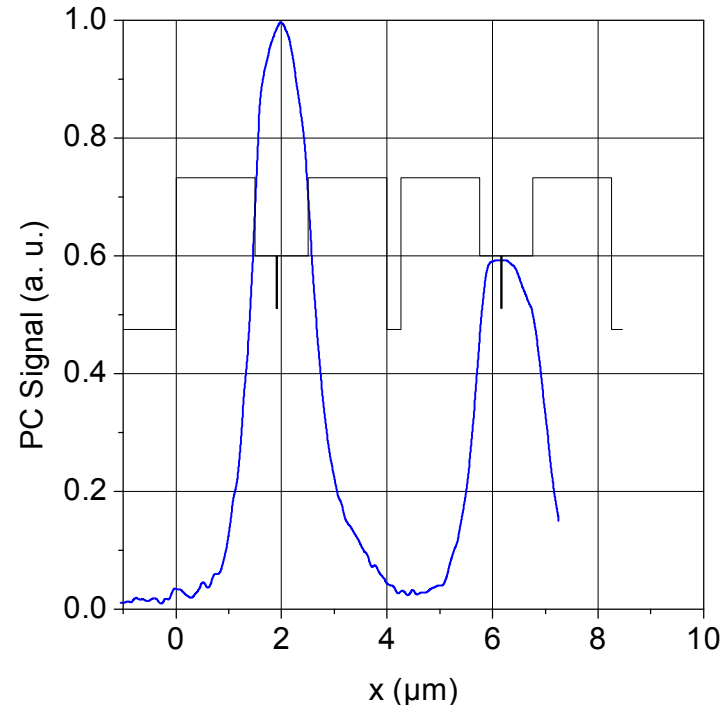
excitation photon energy: 1.69 eV
detection photon energy: 1.5 eV

PL-signal $\sim \delta n$

resonant and non-resonant NSOM Photocurrent (Absorption)



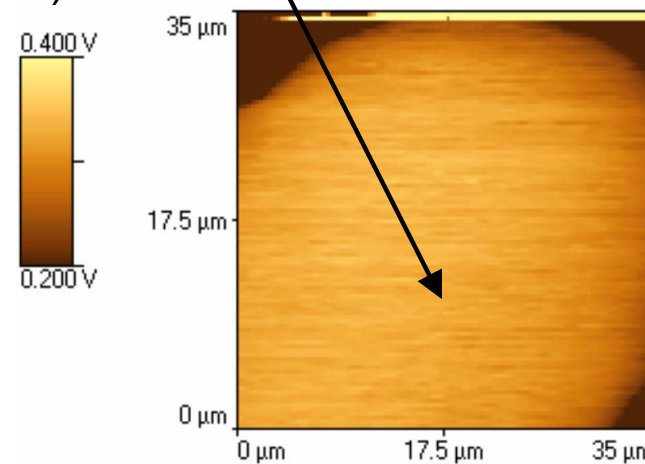
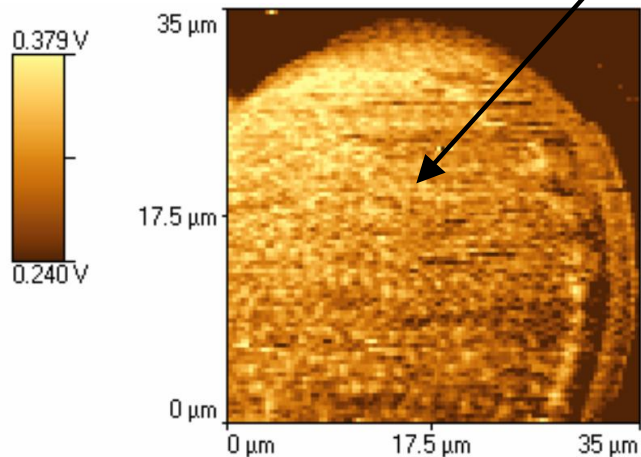
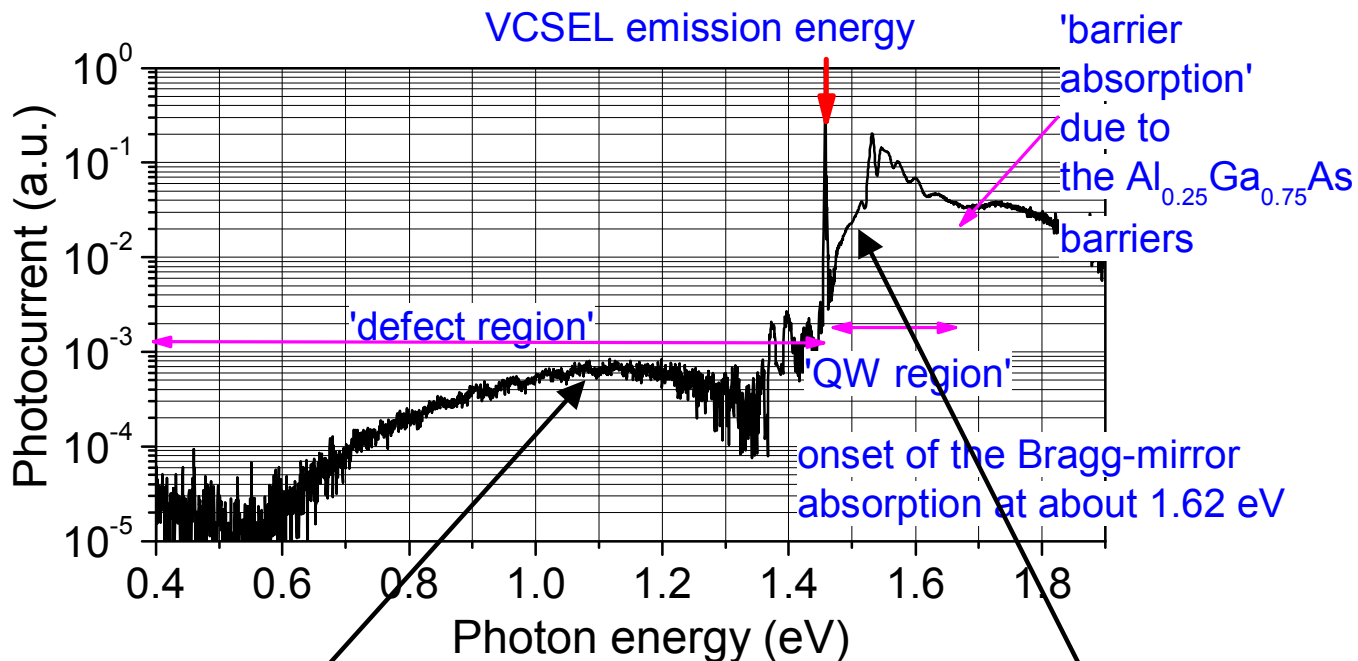
excitation photon energy: 1.69 eV



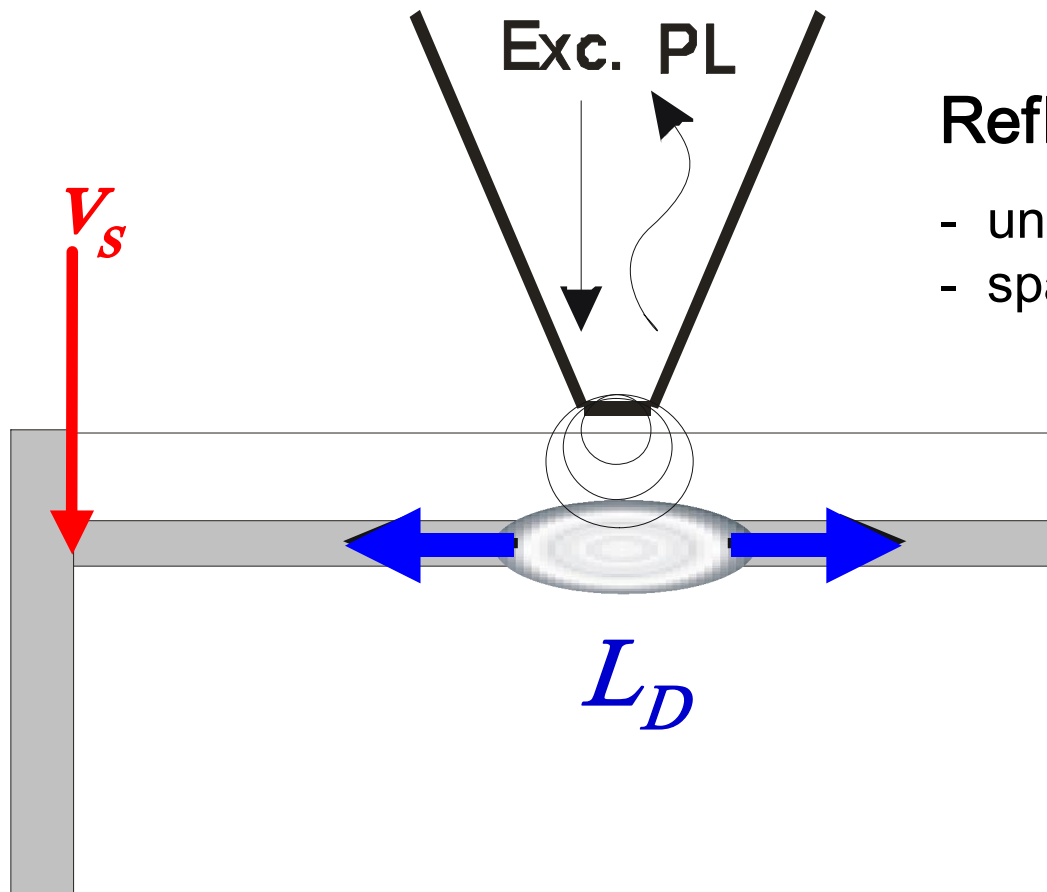
excitation photon energy: 2.8 eV

$$\text{Photocurrent-signal} \sim \delta n \times \text{grad}(V)$$

V. Malyarchuk et al. "Uniformity tests of individual segments of interband cascade diode laser nanostacks[®]" *J. Appl. Phys.* **92**, 2729-2733 (2002).



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- 7. Surface recombination velocity at facets**
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Reflection PL experiment

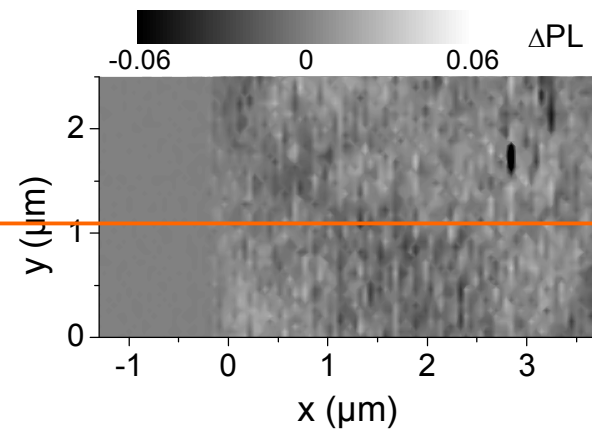
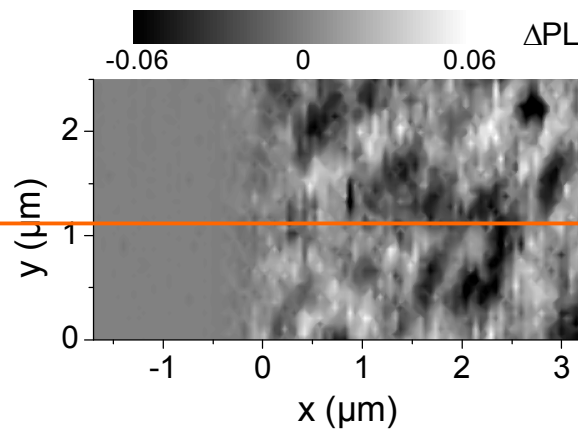
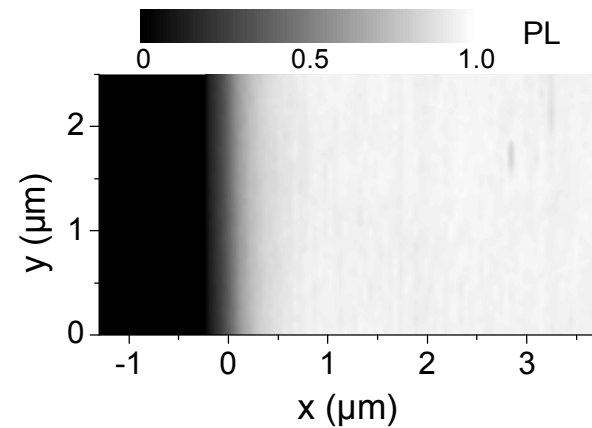
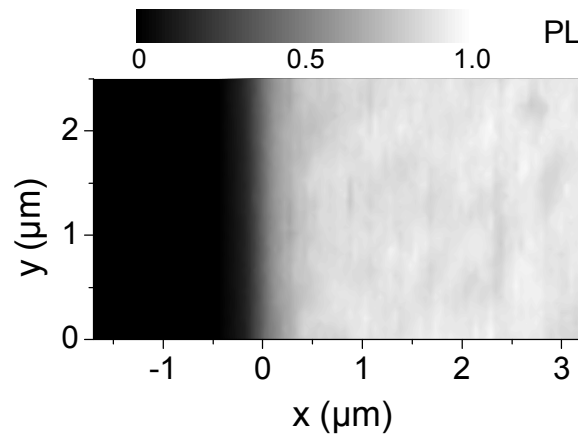
- uncoated fiber tip
- spatial resolution: 150 nm

V. Malyarchuk, J. W. Tomm, V. Talalaev, Ch. Lienau, F. Rinner, and M. Baeumler *Appl. Phys. Lett.* 81 346 (2002).

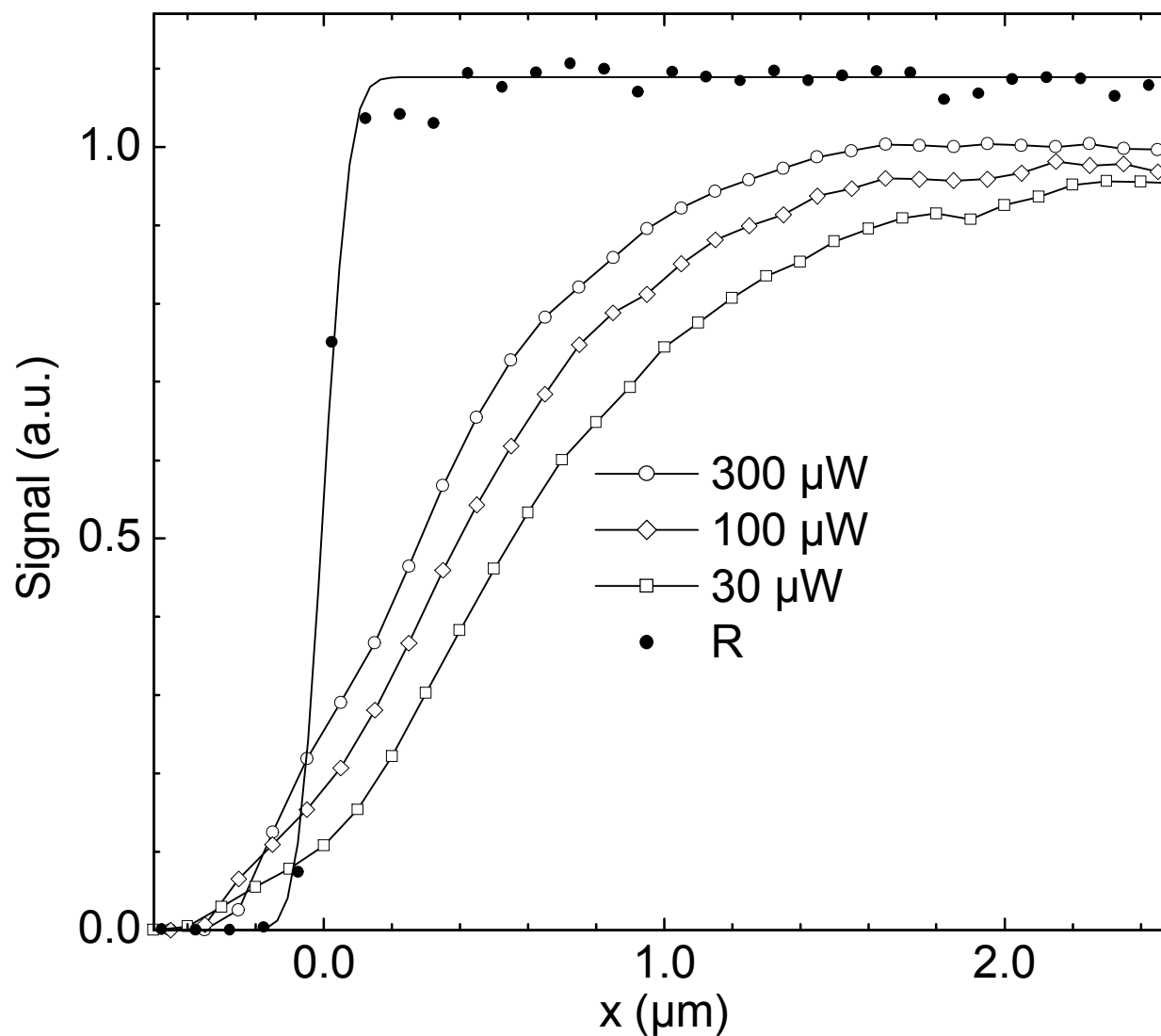
QW PL collected through the fiber and detected by a single photon counting system

$P = 10 \mu\text{W}$

$P = 300 \mu\text{W}$



Line-leveled images



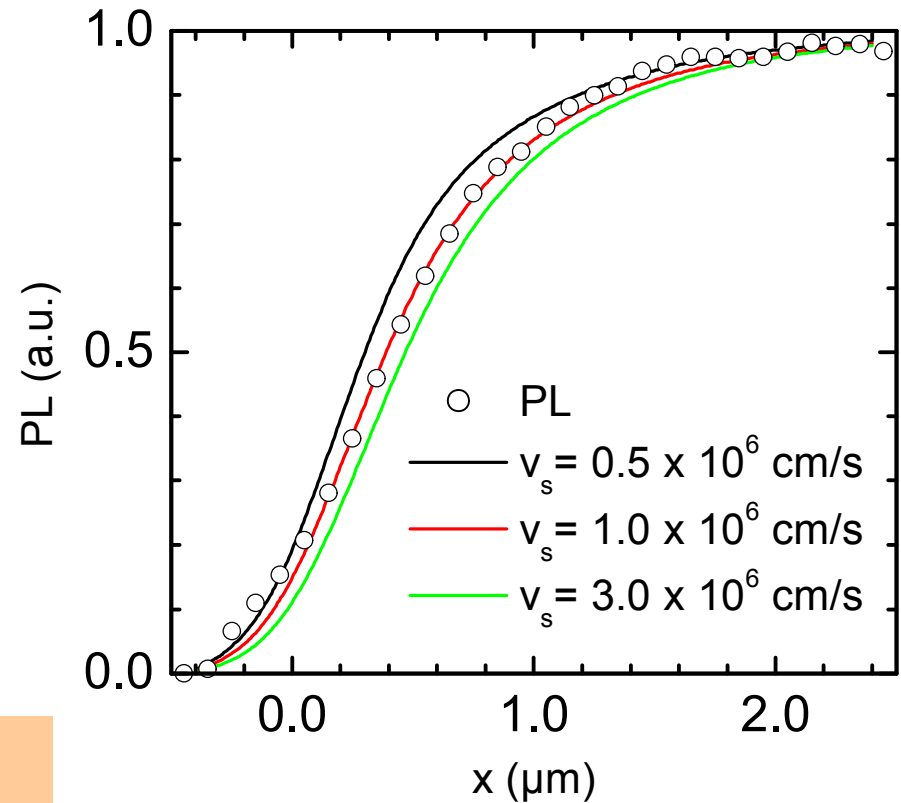
$$D\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right)n - \frac{n}{\tau_{rec}} + n_a g(x_0, y_0) = 0$$

$$L_D = \sqrt{D\tau_{rec}}$$

$$g(x_0, y_0) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(x-x_0)^2 + (y-y_0)^2}{2\sigma^2}}$$

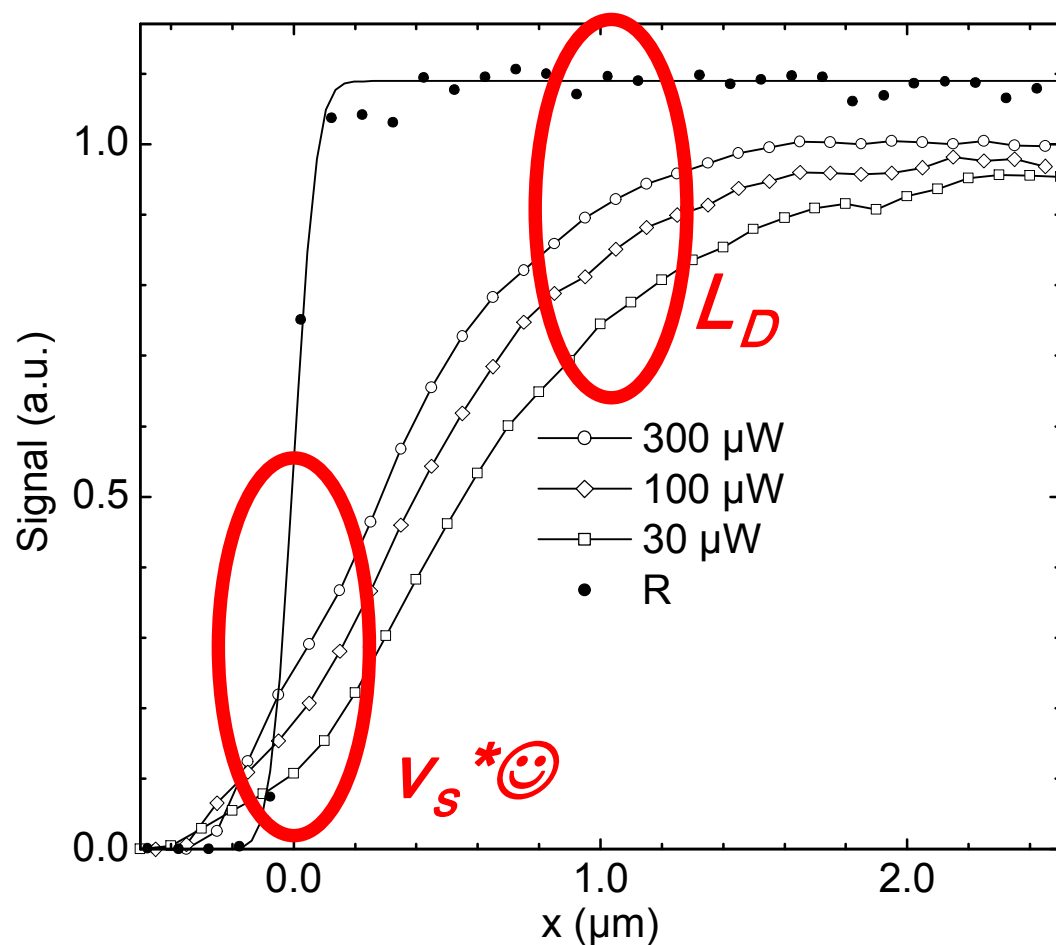
$$x=0: \quad D \frac{\partial n}{\partial x} \Big|_{x=0} = v_s \cdot n$$

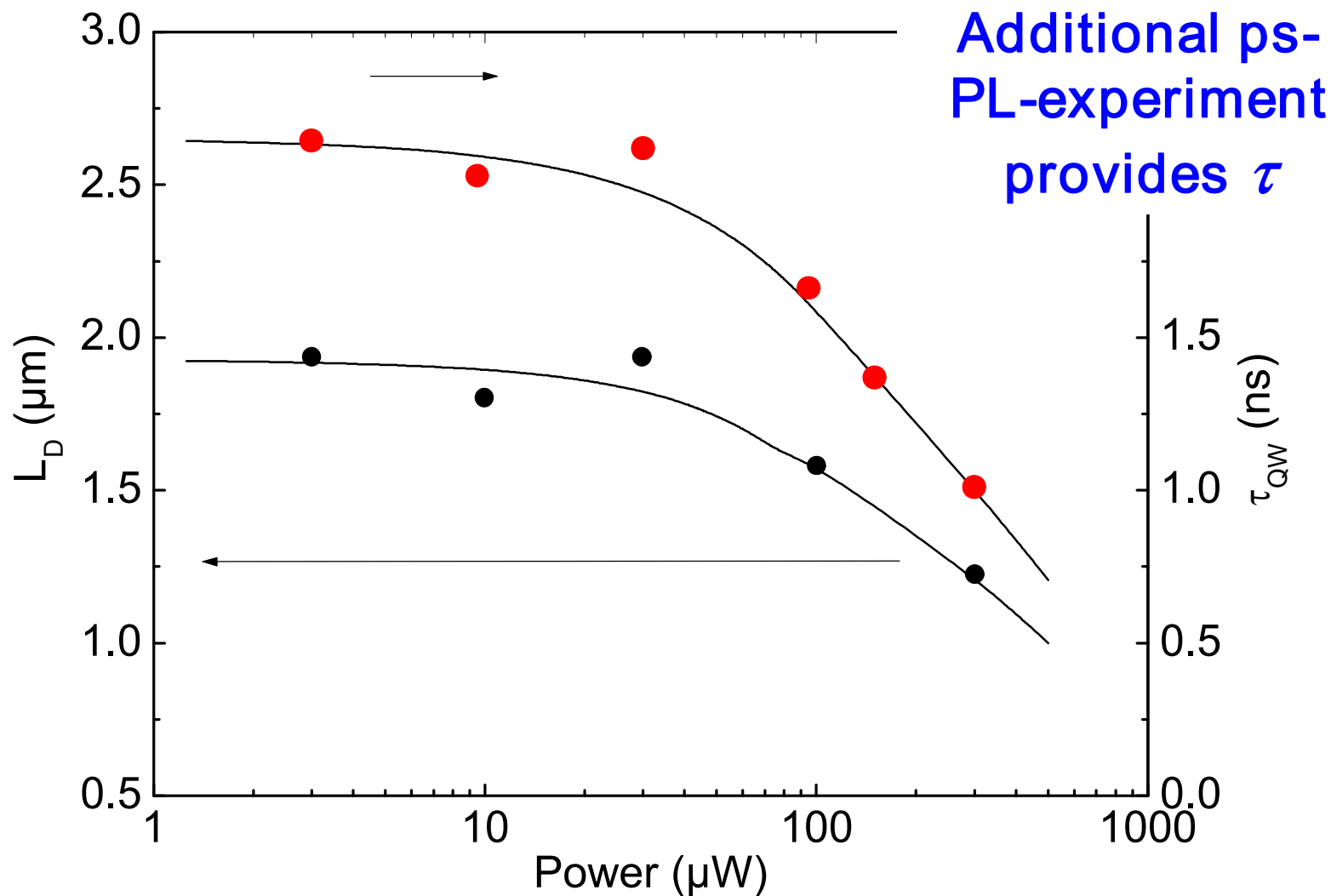
diffusion length L_D and the product $(v_s \cdot \tau_{rec})$ can be deduced

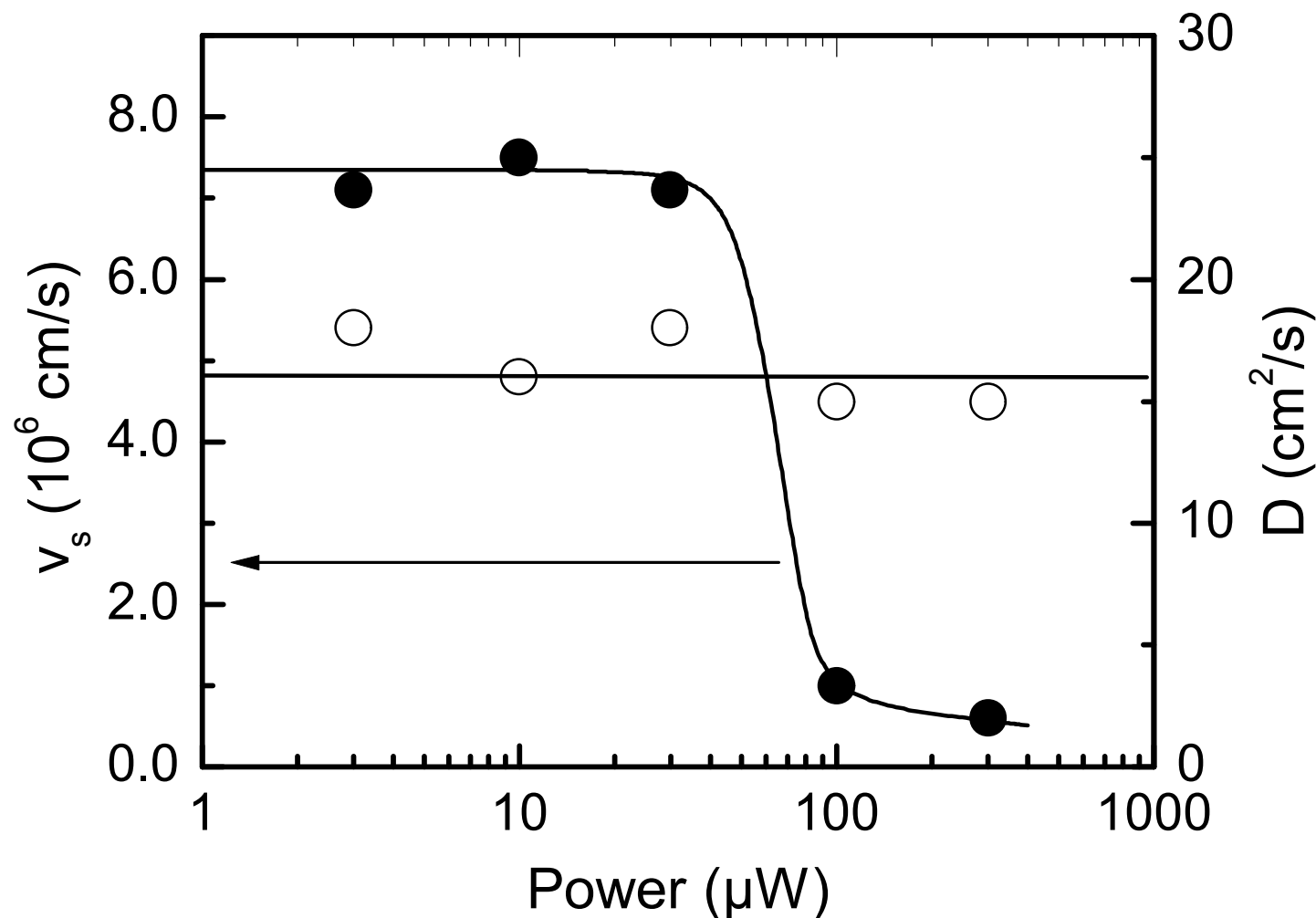


Example calculated for:

$$D = 15 \text{ cm}^2/\text{s} \quad \tau = 1.7 \text{ ns}$$



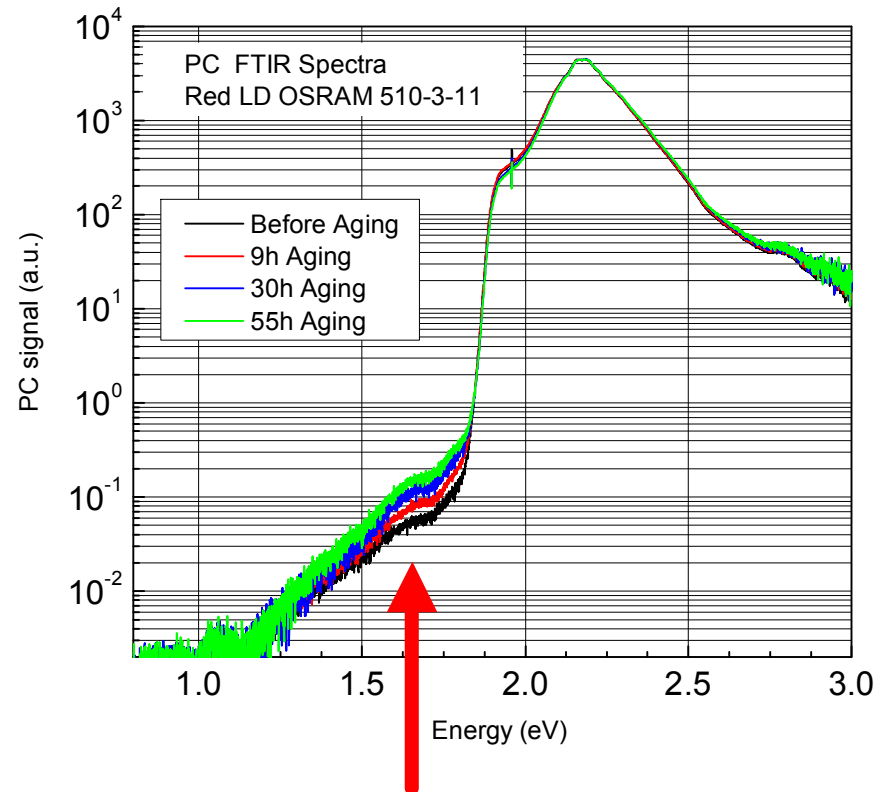
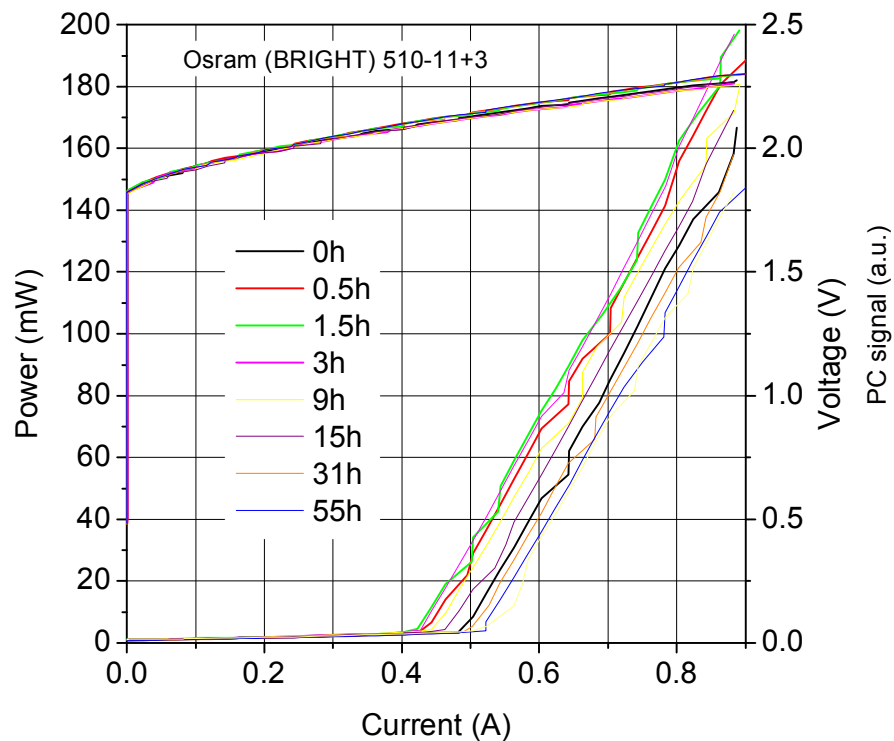


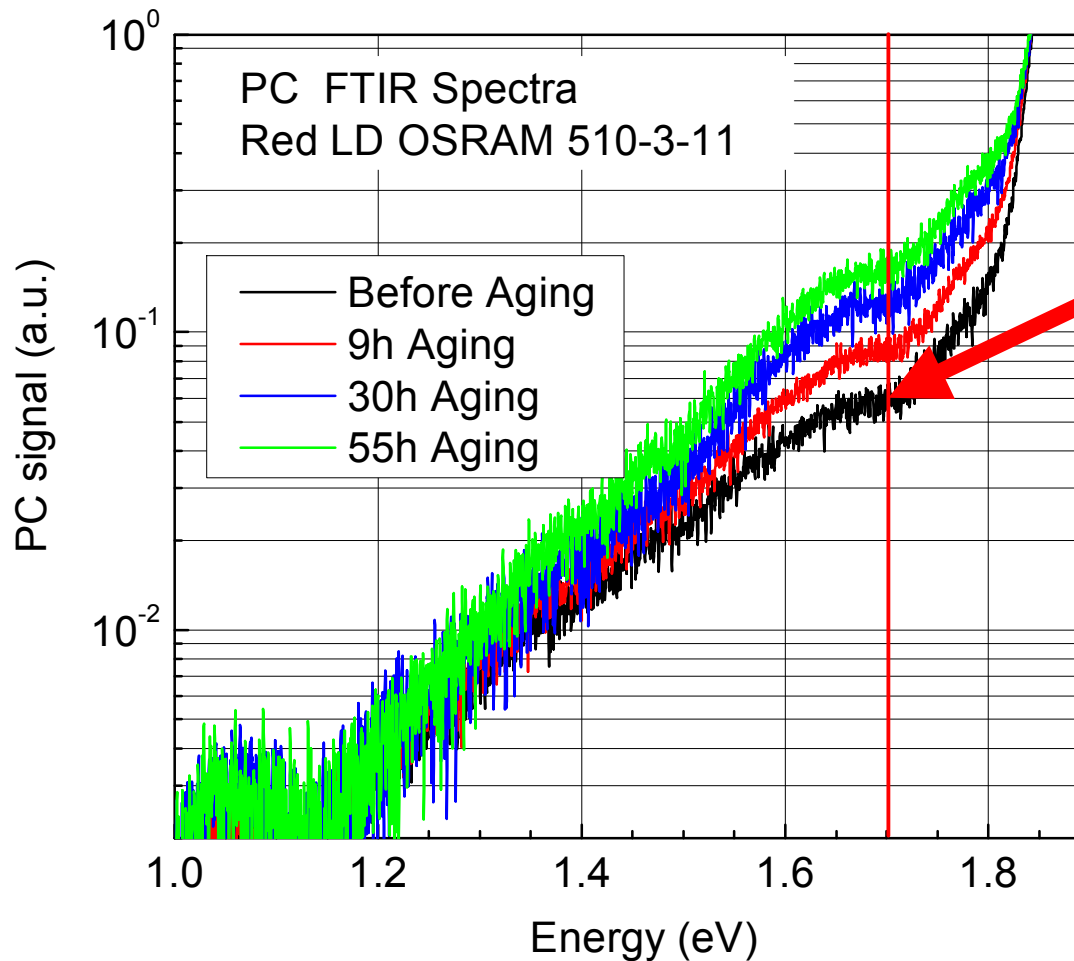


V. Malyarchuk, J. W. Tomm, V. Talalaev, Ch. Lienau, F. Rinner, and M. Baeumler
Appl. Phys. Lett. **81** 346 (2002).

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Result from the BRIGHT project





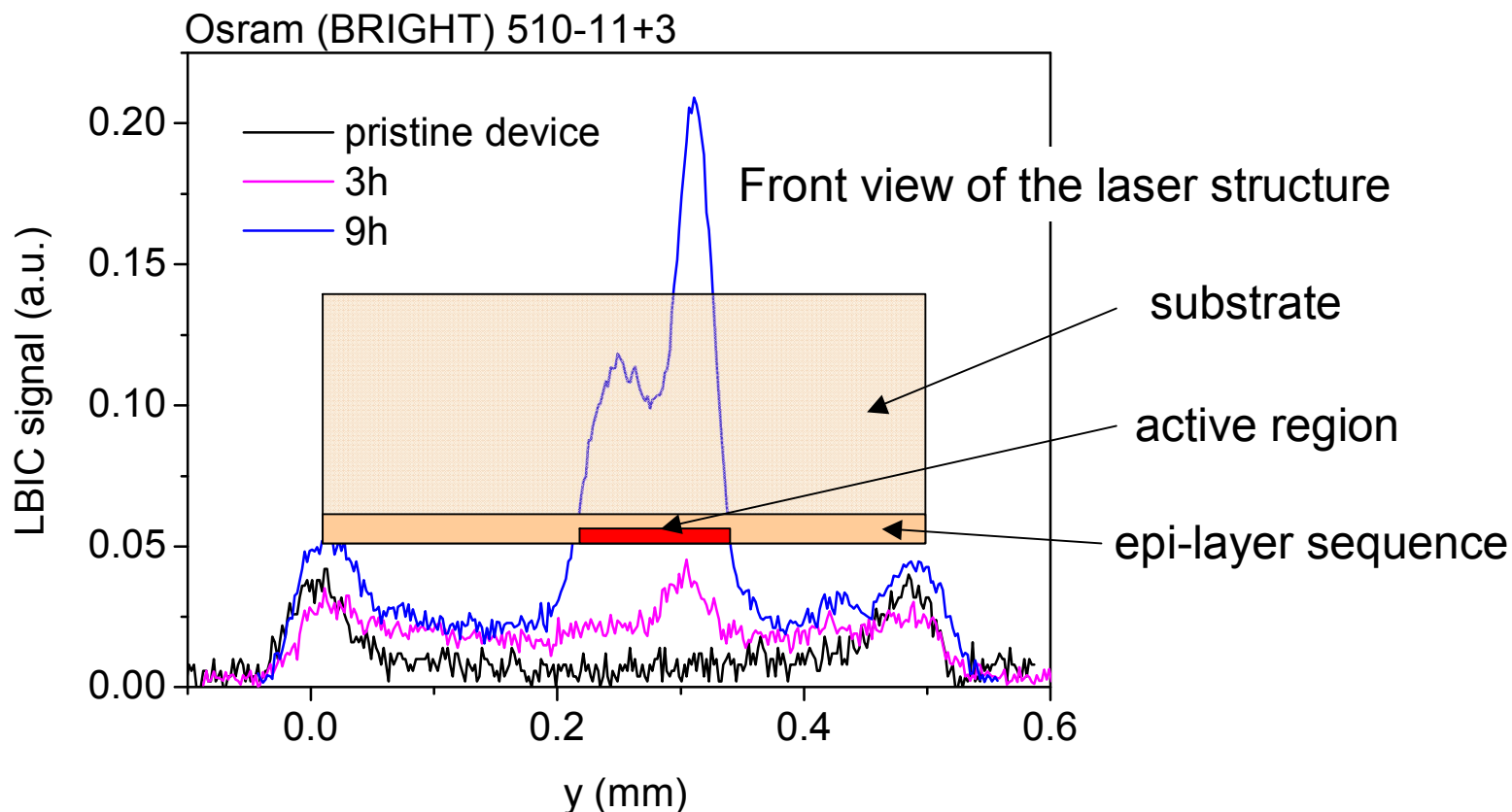
Defect-to-Band-Transitions?

At 1.7 eV (730 nm), there is a threefold increase of the signal within 55 h of aging.

Where are the 'defects' located?

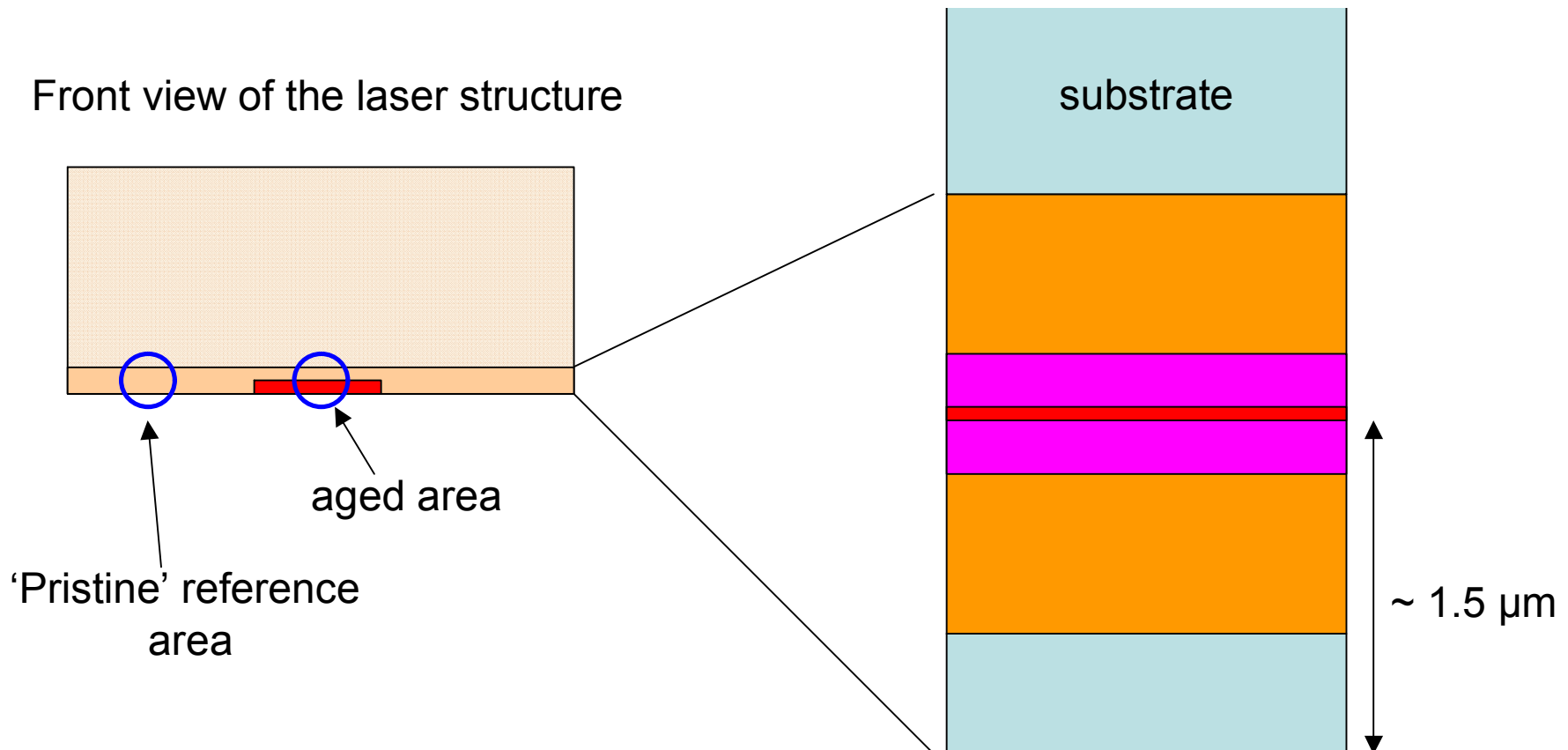
Active region???

1. Laterally, i.e., where along the device...

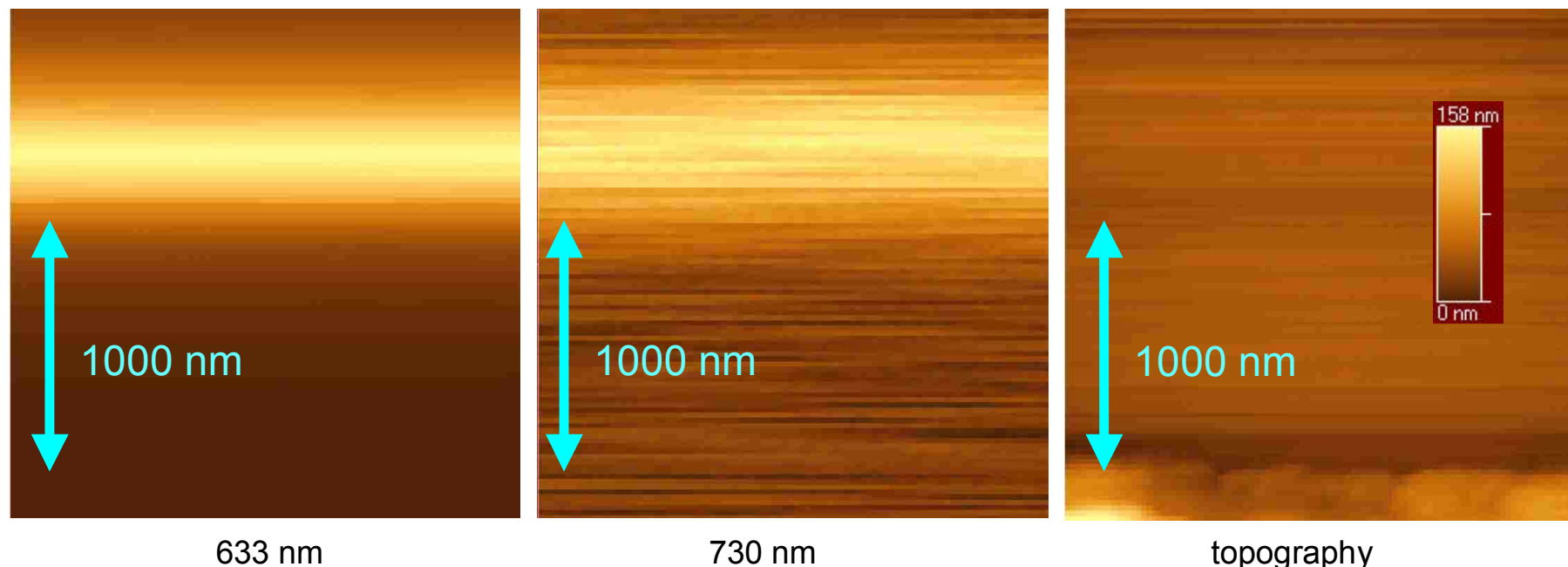


LBIC (Laser Beam Induced Current) excited resonantly to the defects reveals them to be underneath the metalized emitter stripe.

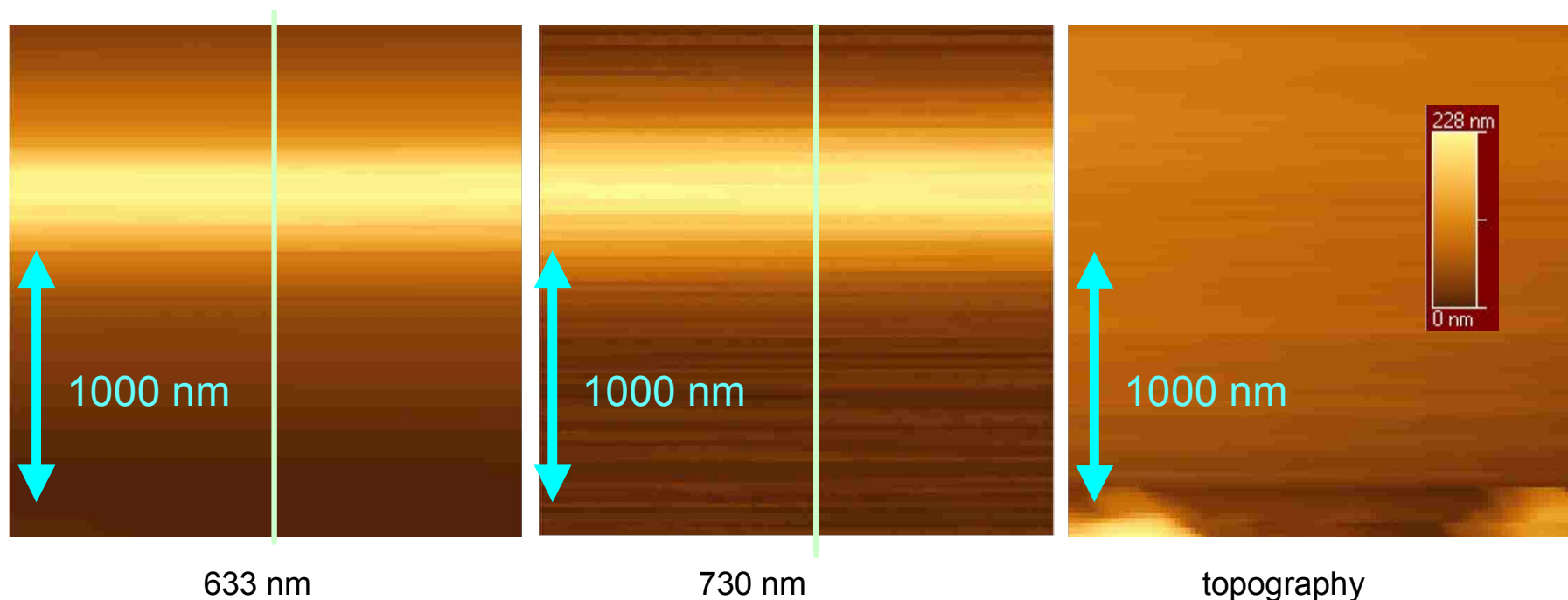
2. Vertically, i.e., where along the layer sequence ...



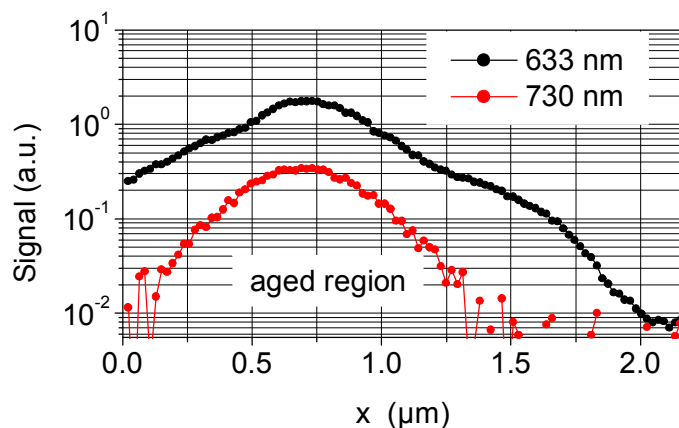
NOBIC-Data: (Nearfield Optical Beam Induced Current)

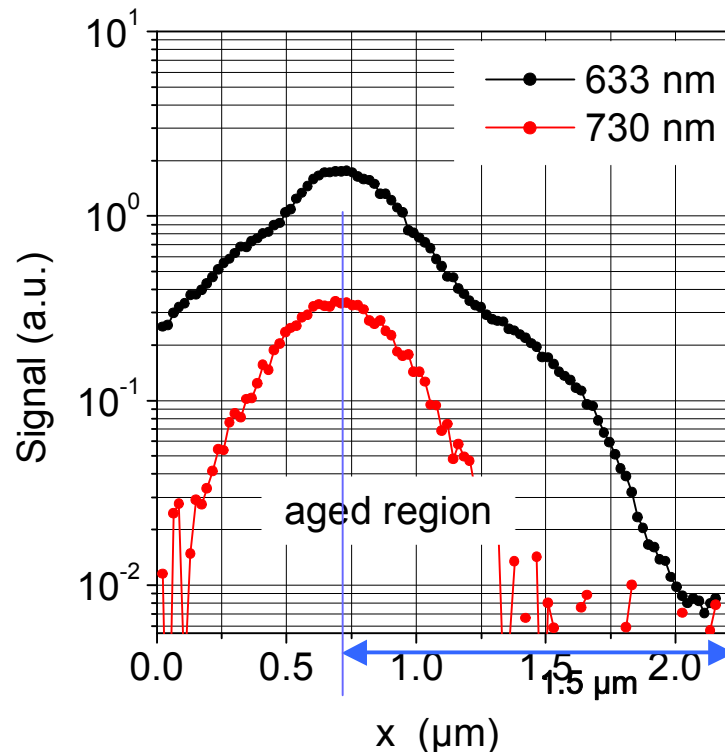
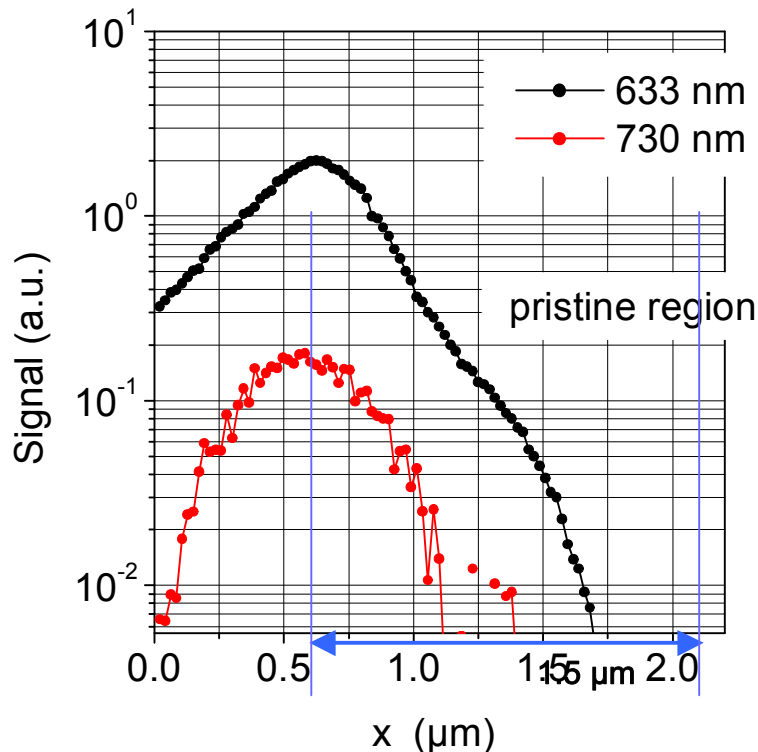


pristine device region



aged device region





There is a threefold increase of the 730 nm- signal with respect to the 633 nm-signal 1.5 μm away from the active region (towards the heat sink), there is no photosensitivity.

The creation of deep levels takes place at a location that allows interaction with the laser emission.

There are additional 3 sets of data couples from different regions, which confirm the result.

Claus Ropers et al. *Appl. Phys. Lett.* **88**, 133513 (2006).

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VCSEL

Surface recombination velocity at facets

Defect creation during device operation

NSOM is extremely useful if you need a spatial resolution beyond the diffraction limit.

NOBIC allows optical analysis at devices along growth direction.

- Christoph Lienau
- Alexander Richter, Tobias Günther, Viktor Malyarchuk, Roland Müller, Claus Ropers, Thomas Elsässer
- Martin Behringer, Johannes Luft, Peter Brick, Norbert Linder, Bernd Mayer, Martin Müller, Sönke Tautz, Wolfgang Schmid, Götz Erbert, Jürgen Sebastian, Siegfried Gramlich, Eberhard Richter, Heiko Kissel, Frank Brunner, Markus Weyers, Günter Tränkle